

Empleando las expresiones del artículo de W.F:Brown

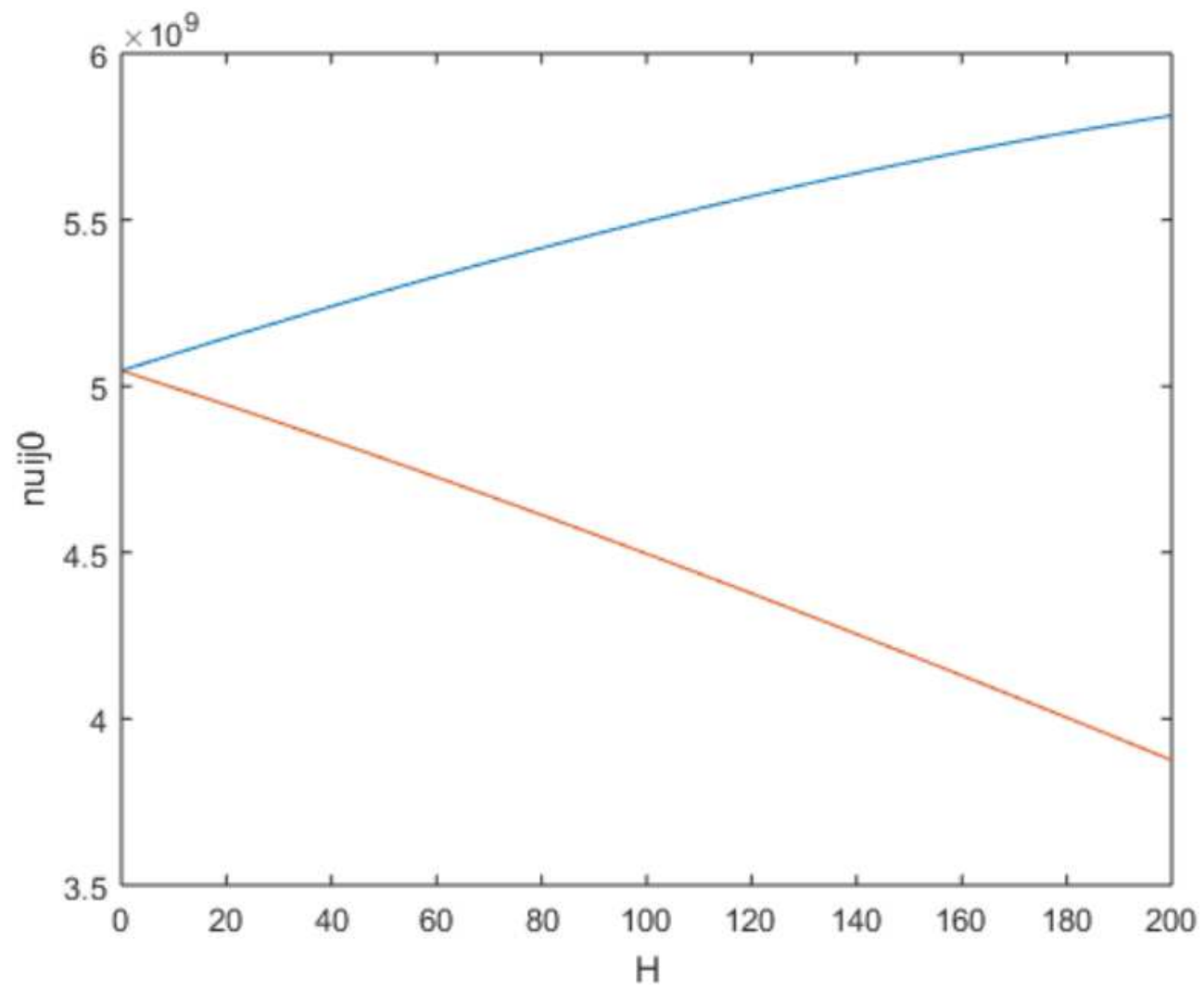
$$\begin{aligned}
 nI_{21} &= b(\beta c'/2\pi)^{1/2} \\
 &\cdot [n_2 c^{(2)} e^{-\beta(V_m - V_2)} - n_1 c^{(1)} e^{-\beta(V_m - V_1)}] \sin \theta_m. \\
 \nu_{ij} &= \nu_{ij}^0 e^{-\beta(V_m - V_i)} \quad (77) \\
 \nu_{ij}^0 &= b(\beta c'/2\pi)^{1/2} c^{(1)} \sin \theta_m. \quad a = \gamma'_0/M_s \quad b = \lambda/M_s^2
 \end{aligned}$$

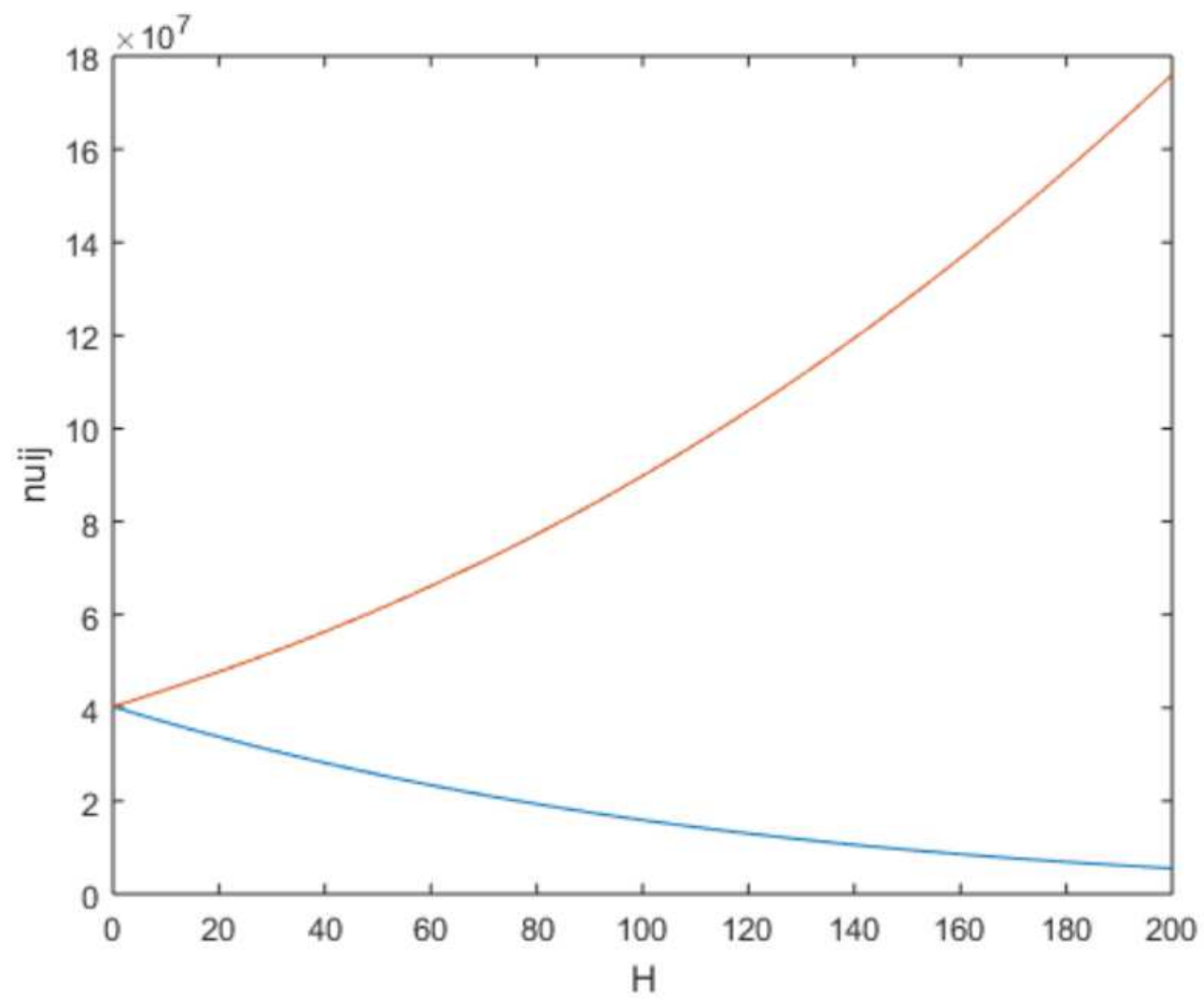
Expresar las ctes. $c^{(1)}$, $c^{(2)}$, c' , θ_m , V_1 , V_2 , V_m , para el caso de anisotropía uniaxial y un capo $H < H_K$ aplicado en la dirección del eje fácil. Usar los valores dados más abajo (repetir para $T = 30$ K)

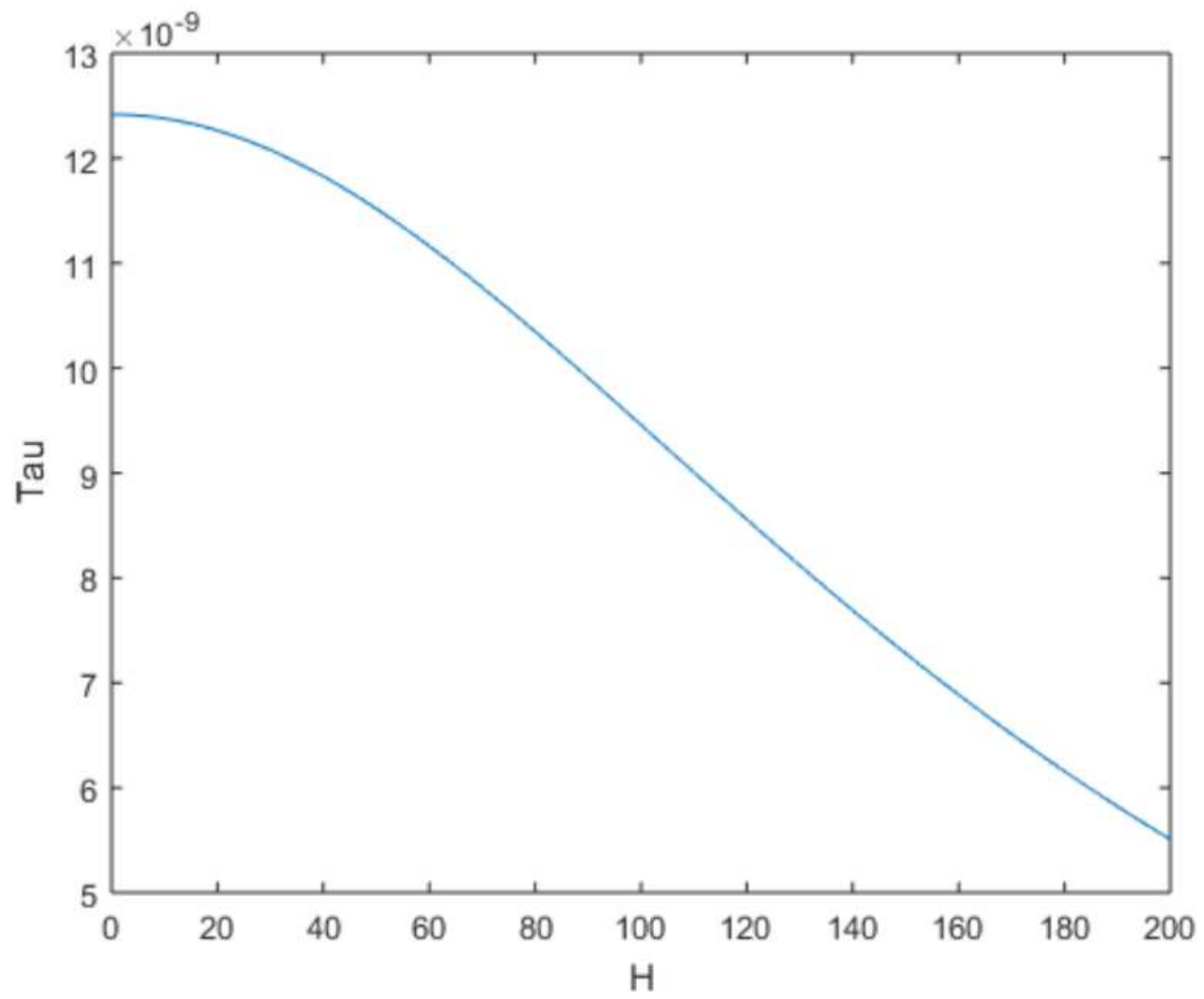
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k=1.38e-16;
Ms=4e2;
lam=10000;
v=1e3*1e-21;
K=2e5;
T=300;
H=[0:1:200];

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Idem para cristales cúbicos

cubic crystal with $K_1 > 0$

$$\nu = G \cdot 2^{1/2} \pi^{-1} b K_1 e^{-(1/4)\beta K_1}$$

$$G = \frac{1}{4} + \frac{1}{4} (9 + 8\rho^2)^{1/2}$$

$$\rho = a/b$$

$$a = \gamma'_0/M_s \quad b = \lambda/M_s^2$$

cubic crystal with $K_1 < 0$,

$$\nu = G \cdot 2^{1/2} (3\pi)^{-1} b |K_1| e^{-(1/12)\beta |K_1|}$$

$$G = -\frac{1}{2} + \frac{1}{2} (9 + 8\rho^2)^{1/2}$$