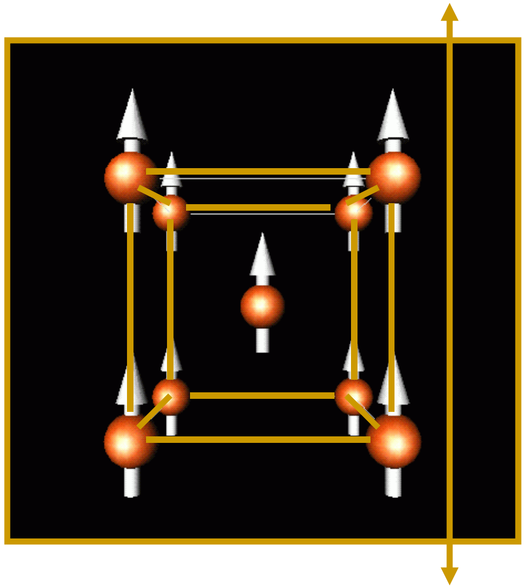
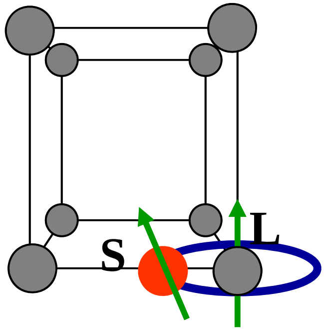
Anisotropía magnetocristalina

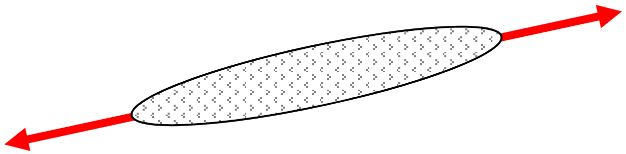
Anisotropía

Fe

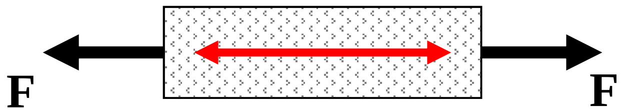




Estructura cristalina → Int. elect - Campo cristalino + interacción Spin – órbita = Anisotropía magneto-cristalina



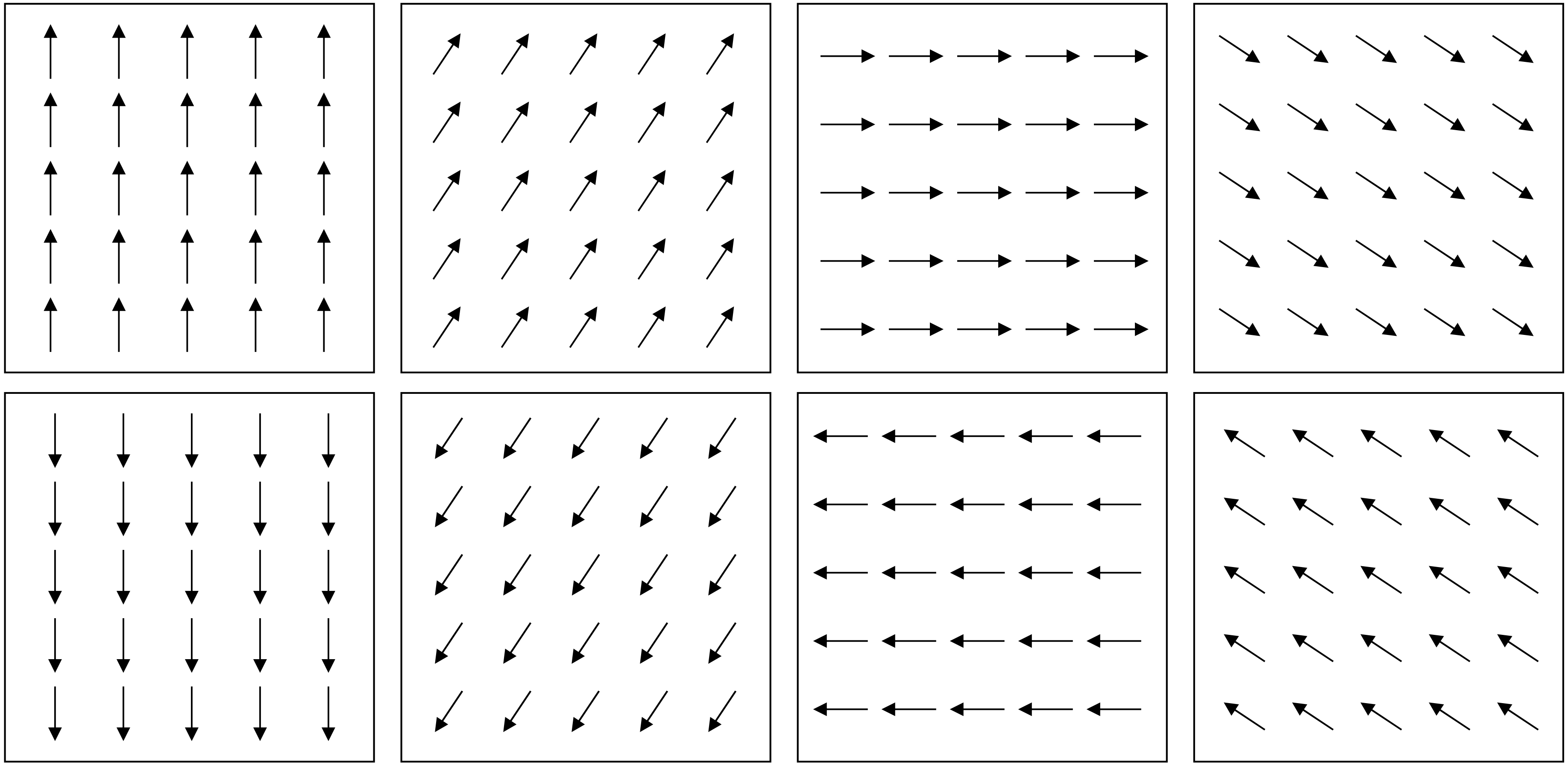
Forma: anisotropía magnetostática



Esfuerzo: anisotropía magnetoelástica

Sólo intercambio
(ausencia de anisotropía)

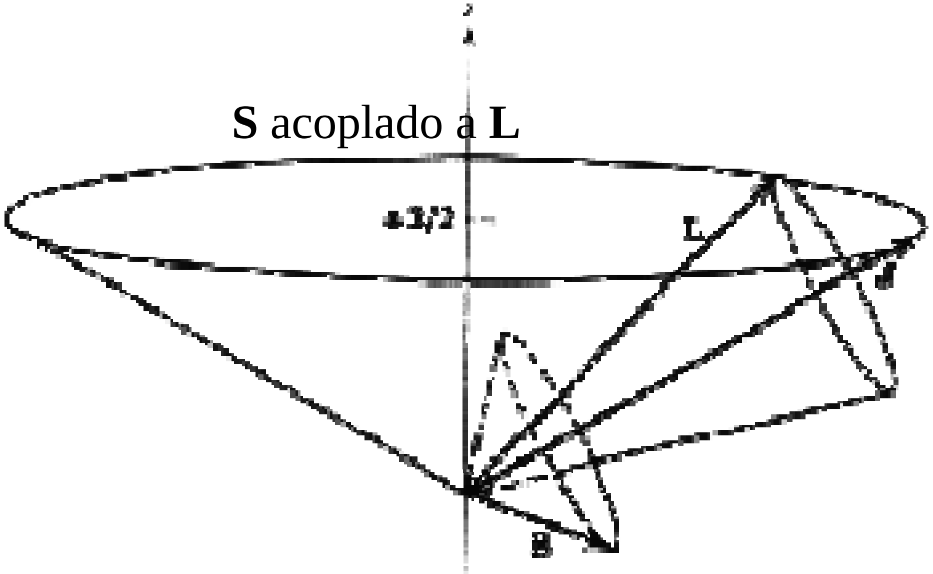
Dirección aleatoria de **M** en 4π => estado continuamente degenerado



Siempre estaríamos en presencia de un superparamagneto

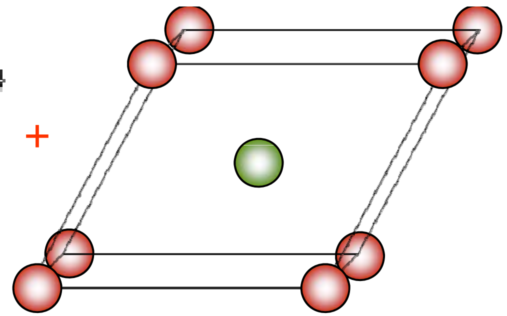
spin – órbita + campo cristalino

S acoplado a **L**

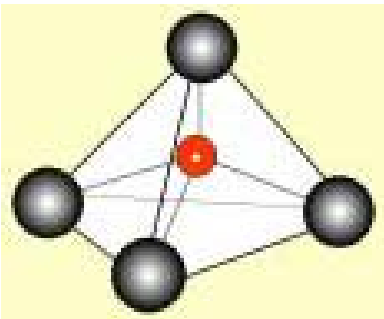


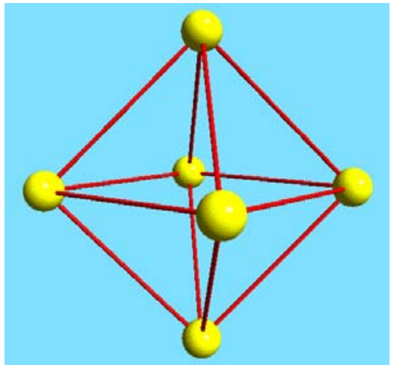
+

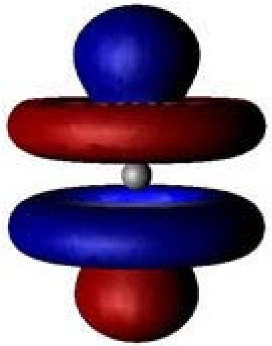
Campo cristalino



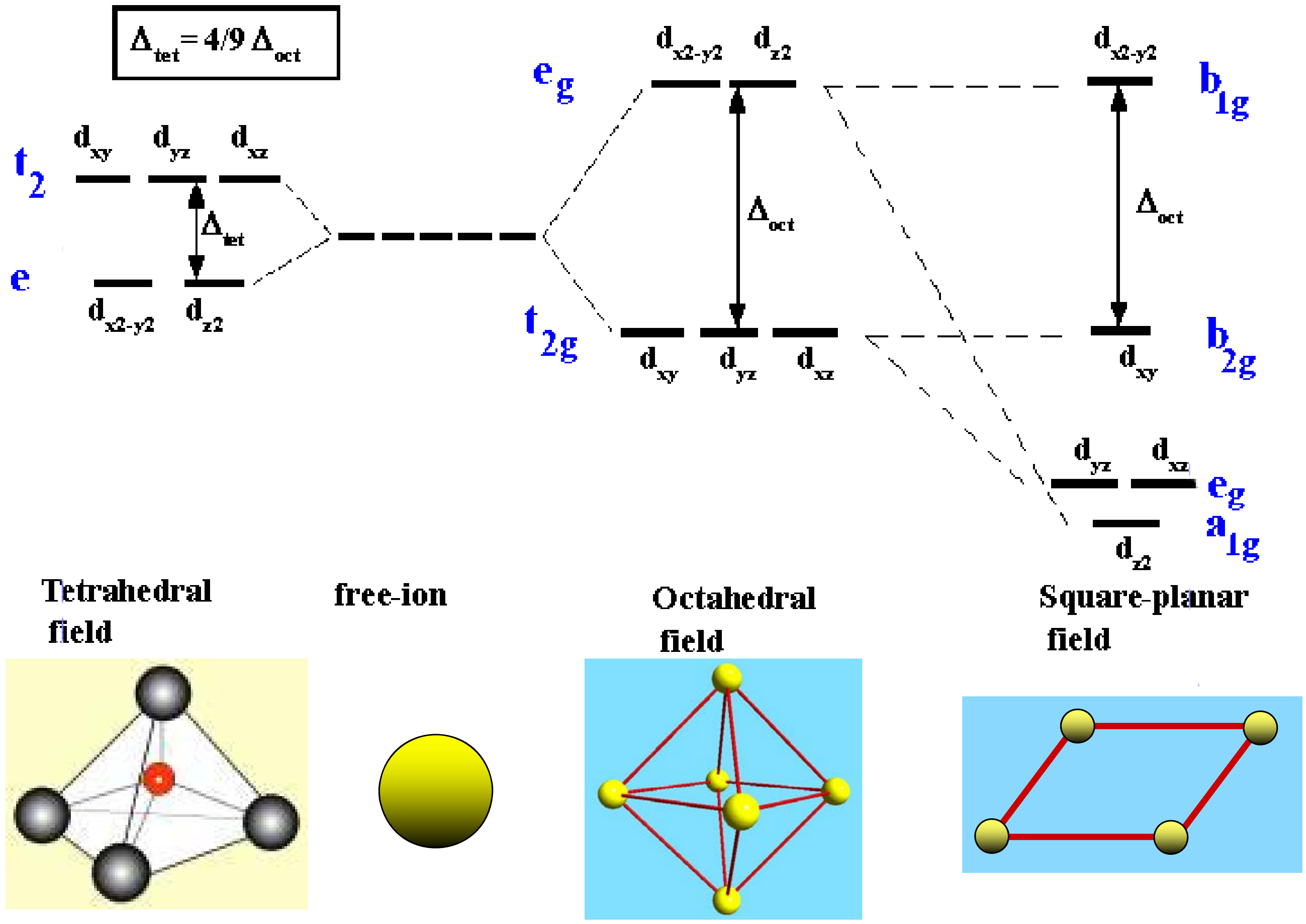
$$= \begin{cases} \textit{anisotropía} \\ \textit{sobre} \\ \vec{L}, \vec{S} \ y \ \vec{J} \end{cases}$$





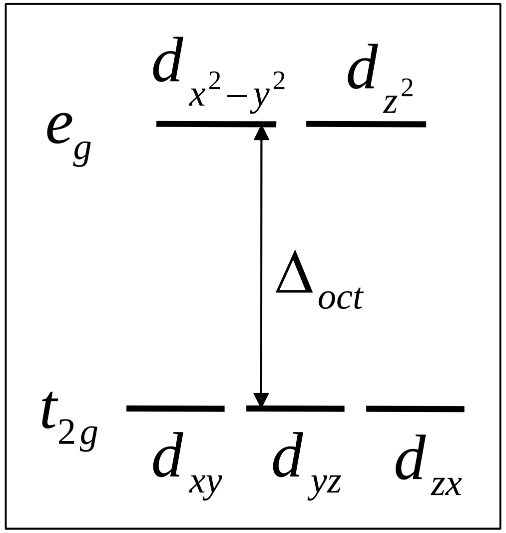


Diagramas de desdoblamiento de orbitales d por el campo cristalino



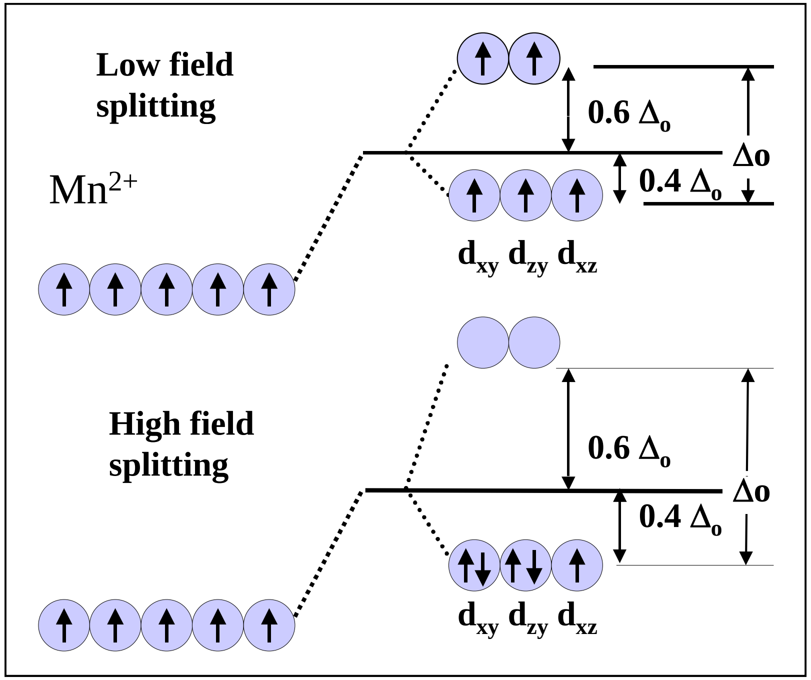
spin – órbita + campo cristalino

Ejemplo 1



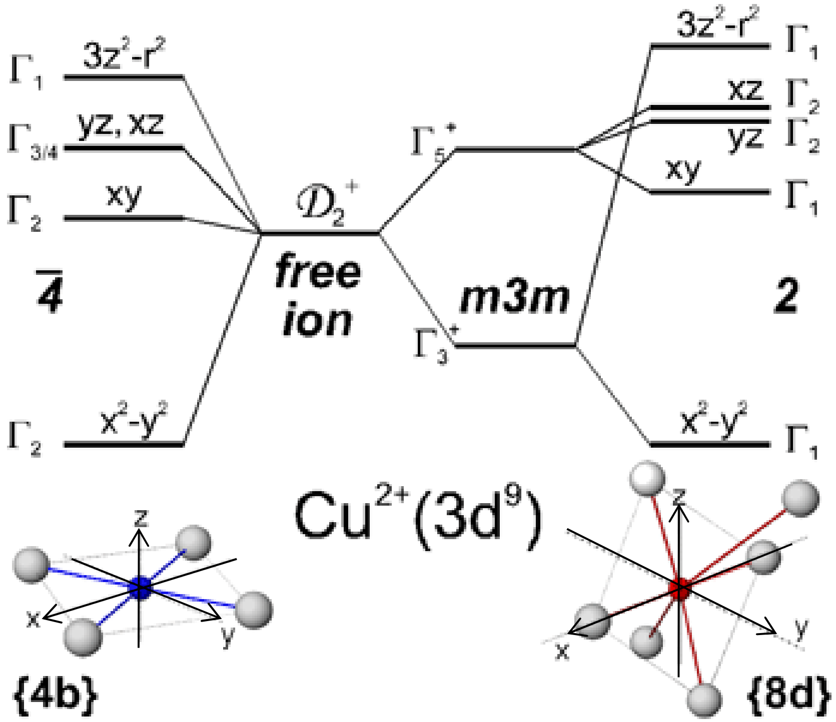
$[Mn^{2+}]=[Ar].3d^5$

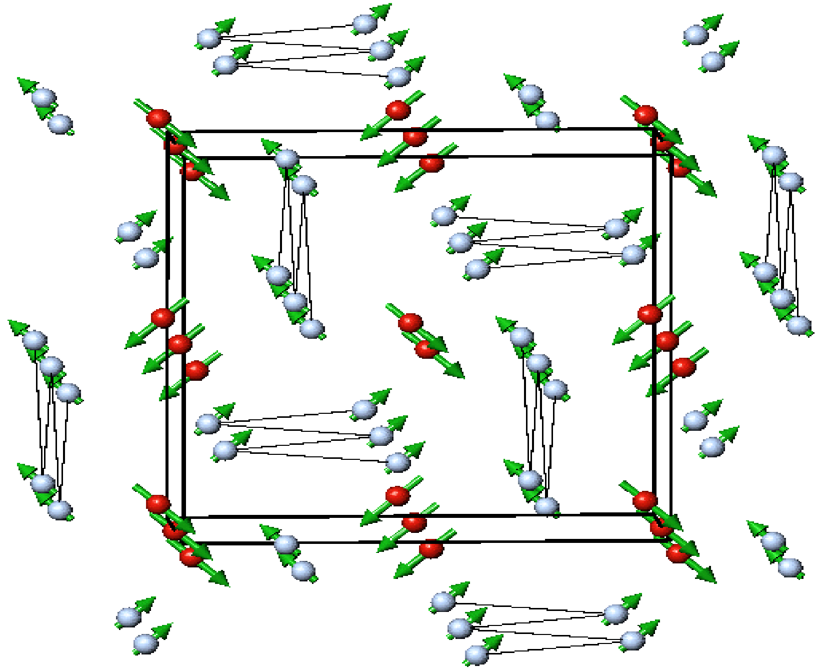
Desdoblamientos bajo campos cristalinos altos y bajos



voh.chem.ucla.edu

Ejemplo 2 Ordenamiento Magnético del CuB_2O_4





Estados electrónicos, simetría local y cordinación de iones Cu^{2+} en sitios 4b y 8d de CuB_2O_4.

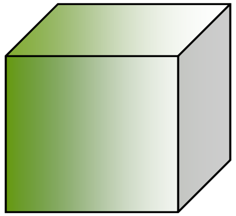
Manfred Fiebig, September 2004 mbi-berlin.de

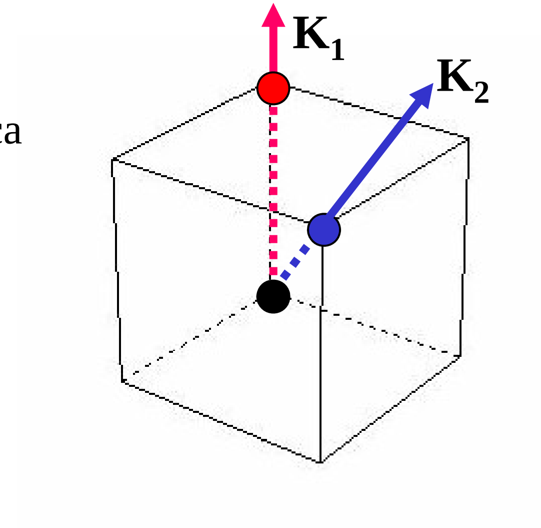
B.Roessli , J.Schefer , G.Petrakovskii, B.Ouladdiaf , M.Böhm , U.Staub, A. Vorotinov and L.Bezmaternikh. Phys. Rev. Letter, (2001).

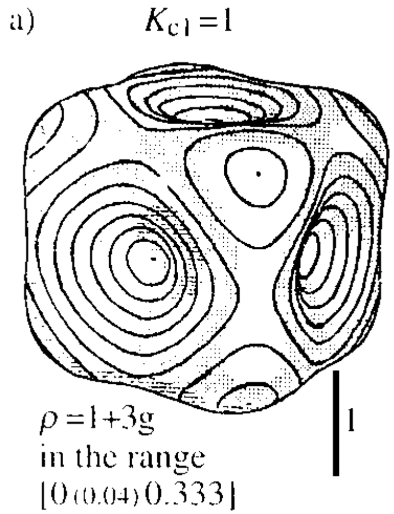
Energía de anisotropía Magnetocristalina

Definiciones

Geometría cúbica

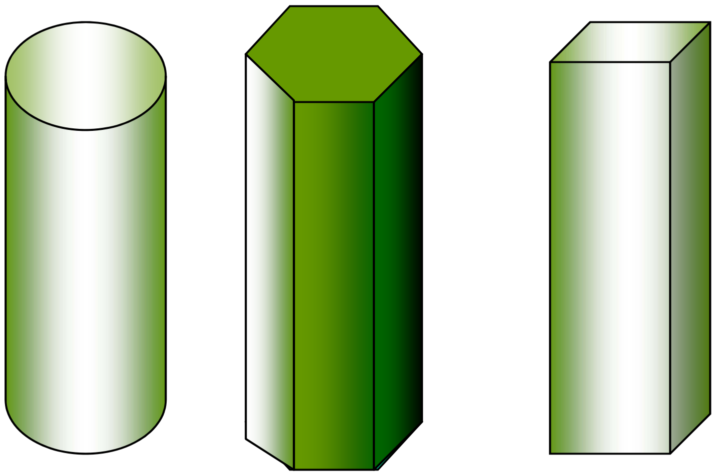


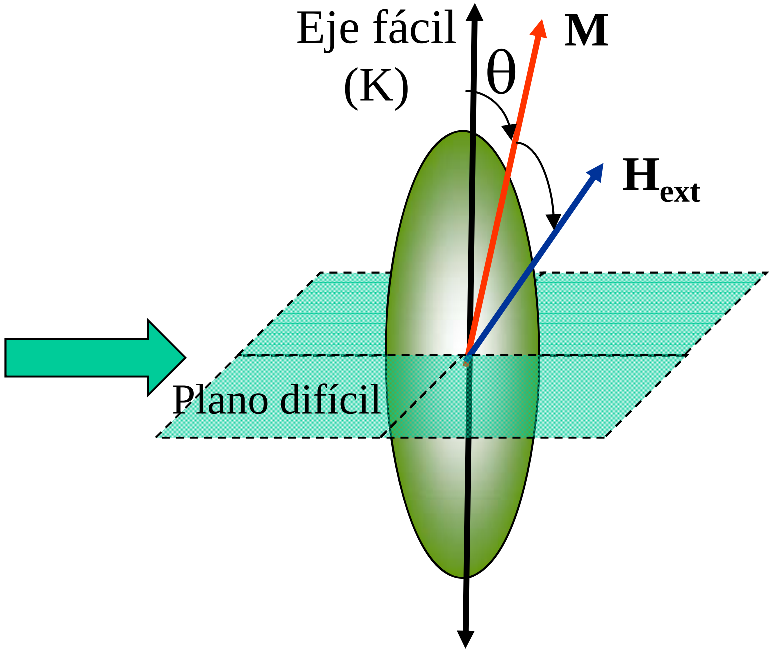




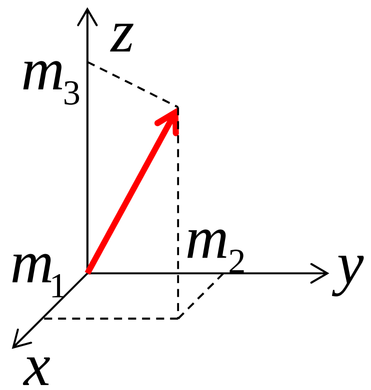
Curvas de energía constante

Geometría cilíndrica, tetragonal, exagonal





Anisotropía – descripción fenomenológica



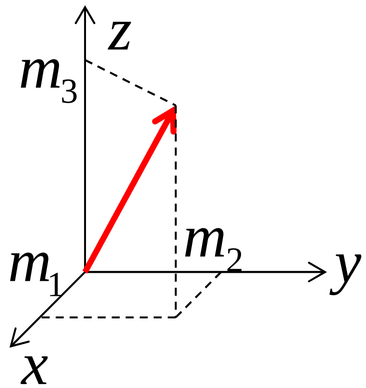
m_i cosenos directores de la magnetización

$$m_{1,2,3} = \frac{M_{x,y,z}}{M} \ldots$$

e_K energía de anisotropía por unidad de volumen

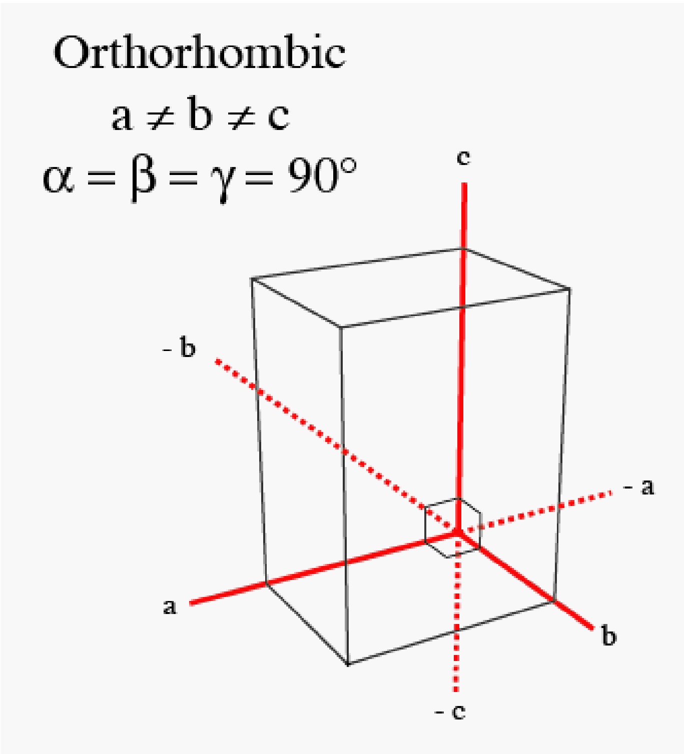
$$e_K = \sum_i K_i m_i^2 + \sum_{ij} K_{ij} m_i^2 m_j^2 + K_{123} m_1^2 m_2^2 m_3^2 + \sum_i K_i m_i^4 + \cdots$$

E_K energía de anisotropía $\quad E_K = \int e_K dV$

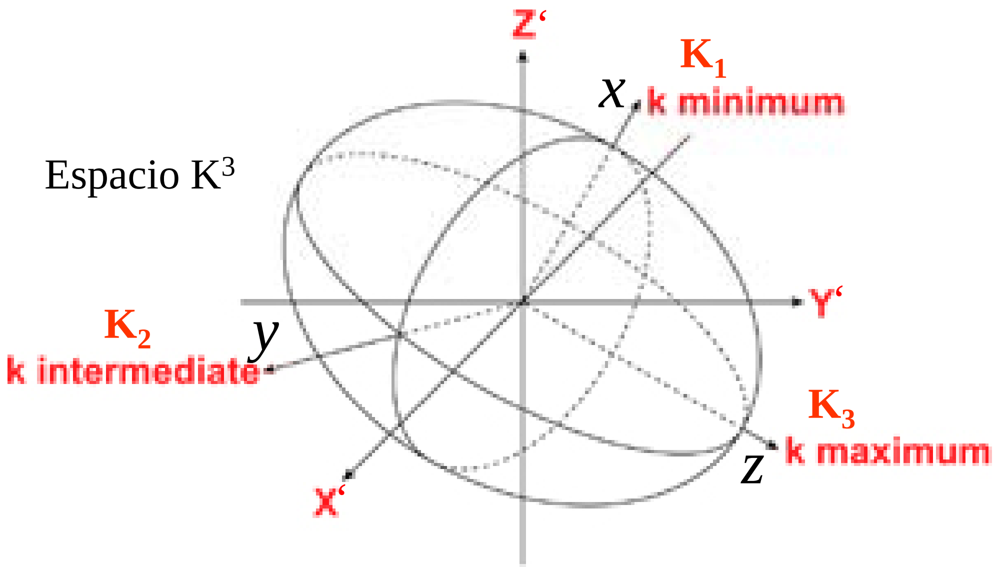


$$e_K = \sum_i K_i m_i^2 + \sum_{ij} K_{ij} m_i^2 m_j^2 + K_{123} m_1^2 m_2^2 m_3^2 + \sum_i K_i m_i^4 + \cdots$$

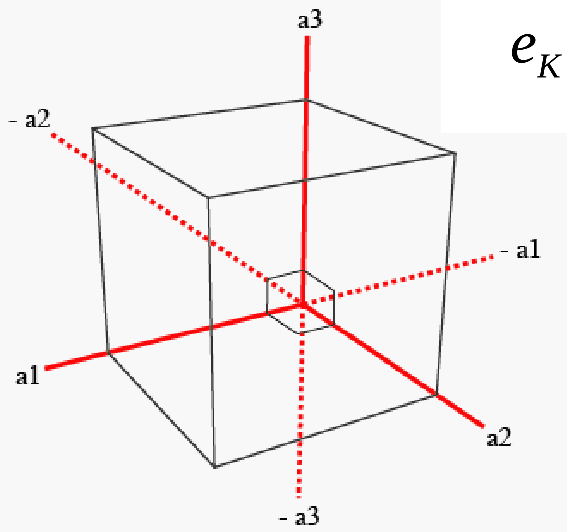
Ejemplo: sistema ortorrómbico



$$e_K \approx K_1 m_1^2 + K_2 m_2^2 + K_3 m_3^2$$



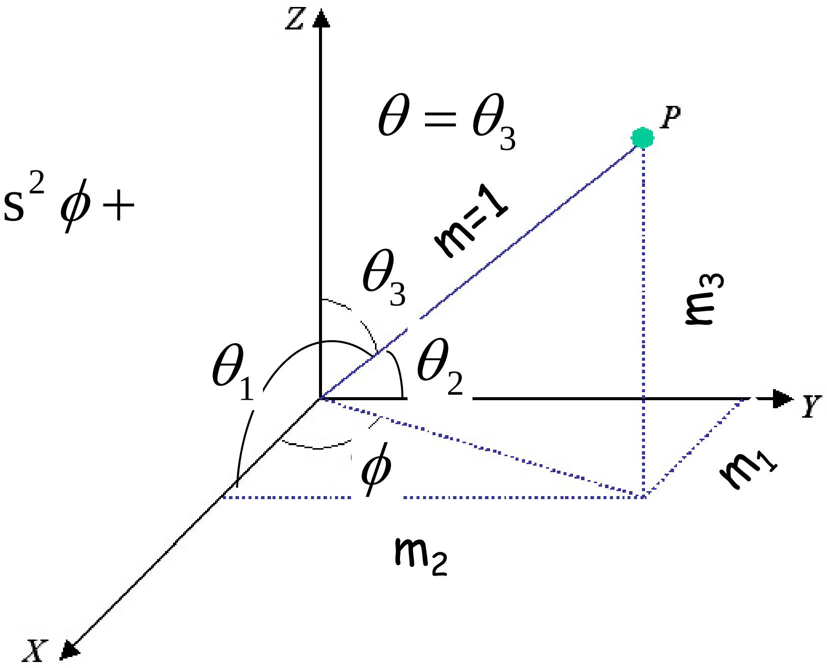
sistema cúbico



$$e_K = K_1\left(m_1^2 m_2^2 + m_1^2 m_3^2 + m_2^2 m_3^2\right) + K_2 m_1^2 m_2^2 m_3^2$$

$$m_i = \cos\theta_i$$

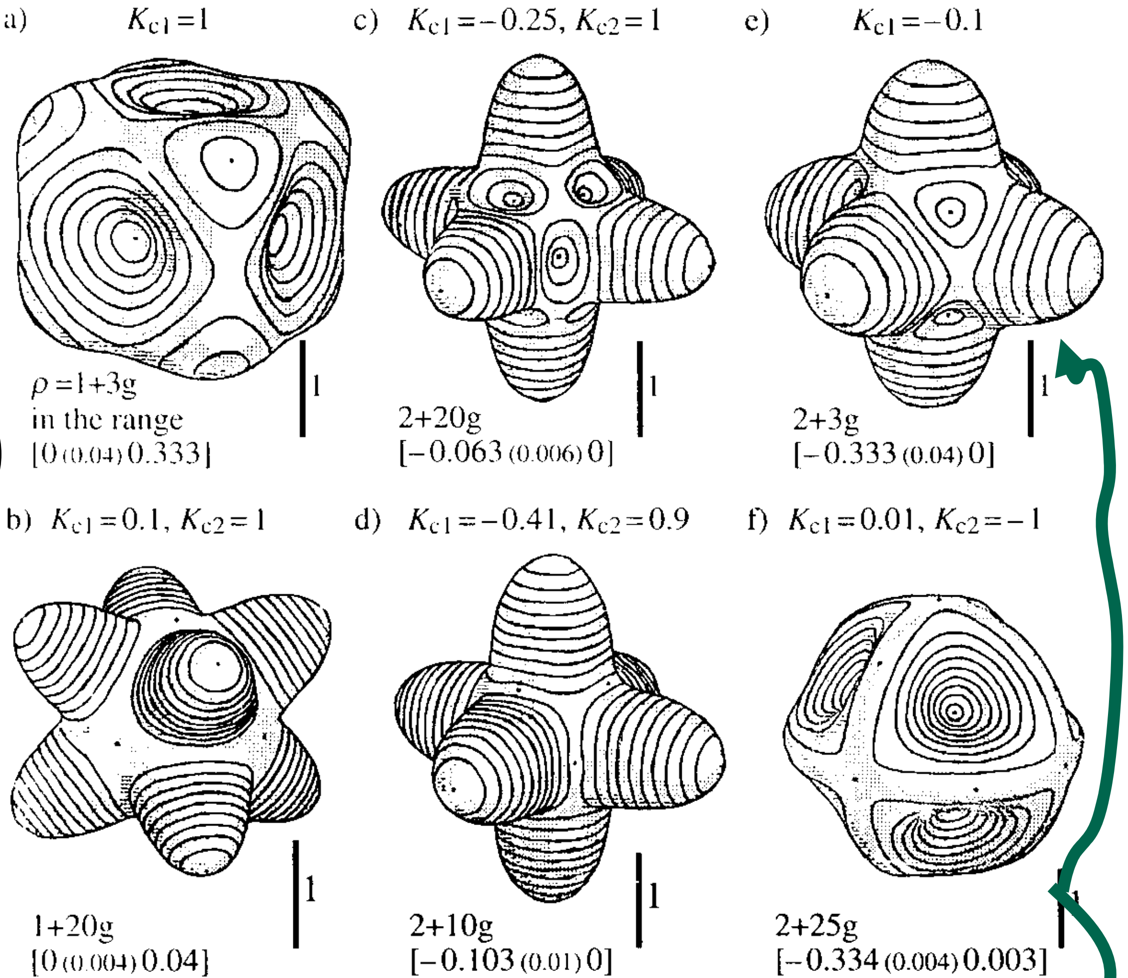
$$e_K = \left(K_1 + K_2\cos^2\theta\right)\sin^4\theta\sin^2\phi\cos^2\phi +$$
$$+ K_1\sin^2\theta\cos^2\theta$$



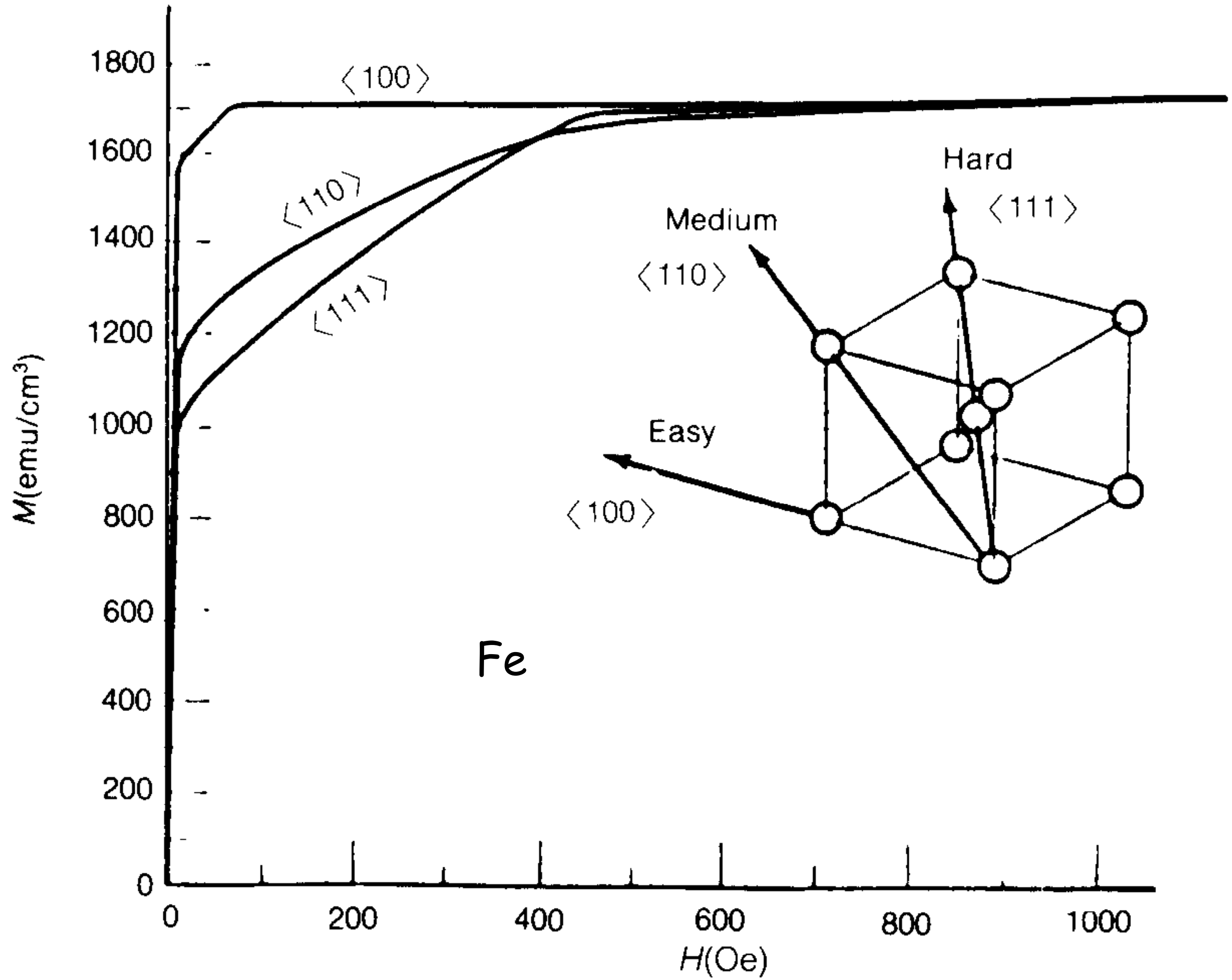
sistema cúbico

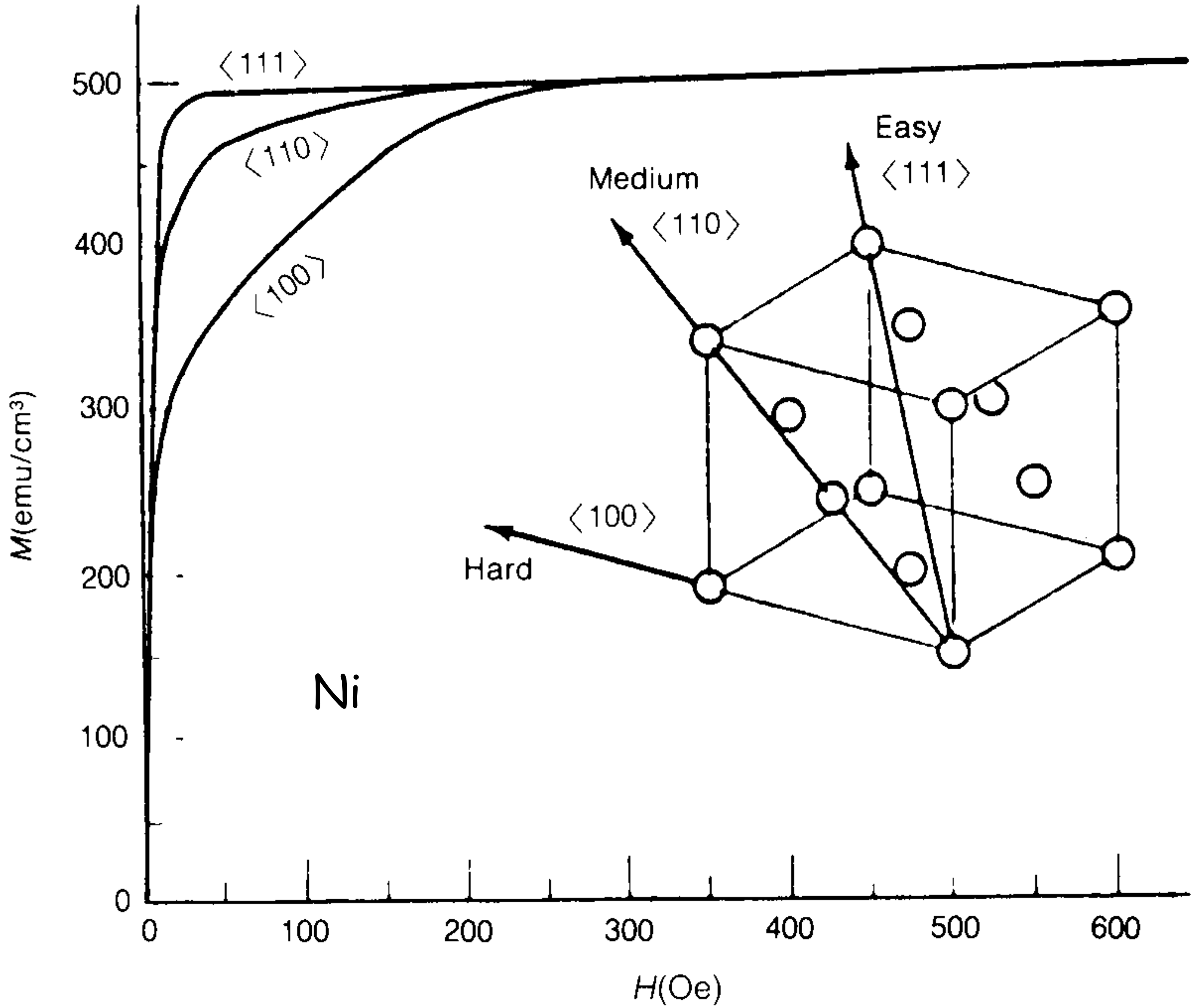
Curvas de energía constante

$$e_K = K_1\left(m_1^2 m_2^2 + m_1^2 m_3^2 + m_2^2 m_3^2\right) + K_2 m_1^2 m_2^2 m_3^2$$



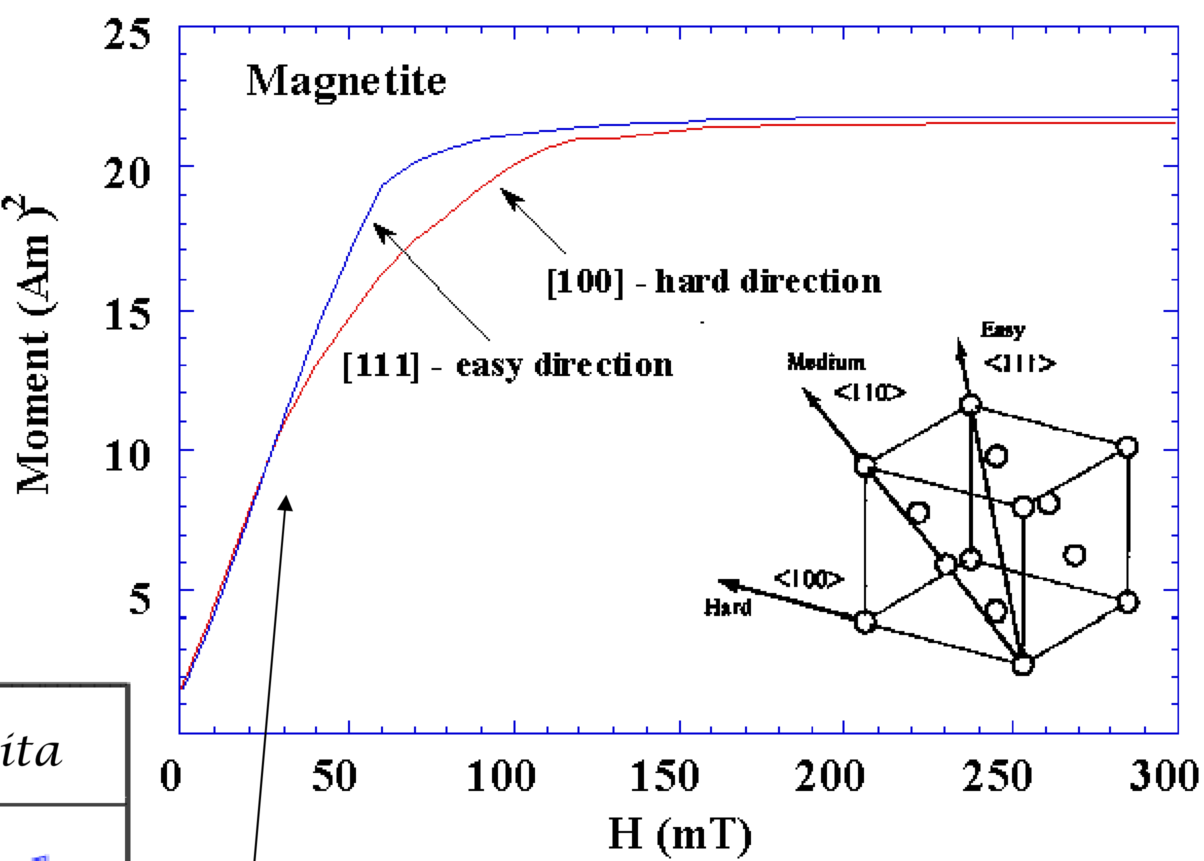
Material	K_1 (10^5 J/m^3)	K_2 (10^5 J/m^3)	Eje fácil
Fe	0.480	0.05	(100)
Ni	-0.045	-0.023	(111)

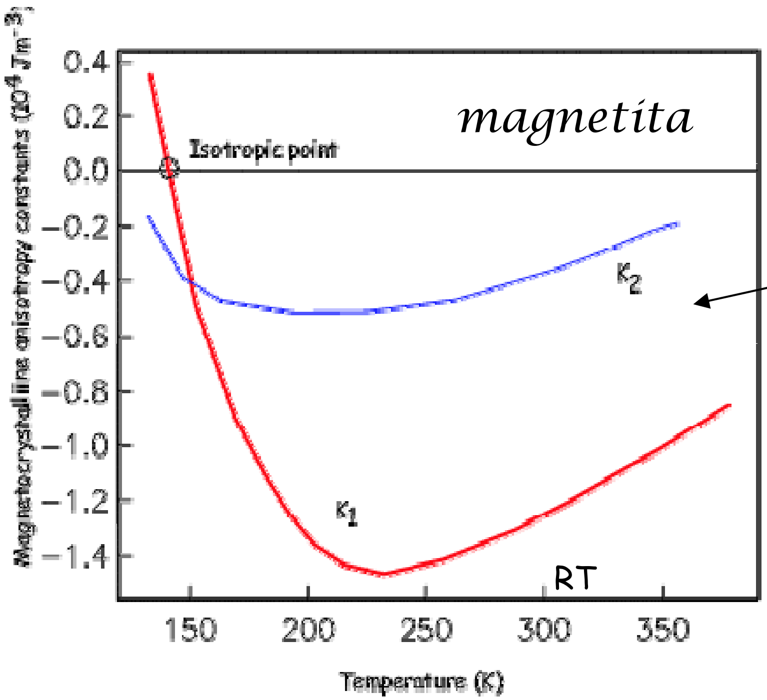

1800
1600
1400
1200
1000
800
600
400
200
0
M(emu/cm³)
0
200
400
600
800
1000
H(Oe)
⟨100⟩
⟨110⟩
⟨111⟩
Hard
⟨111⟩
Medium
⟨110⟩
Easy
⟨100⟩
Fe



⟨111⟩
⟨110⟩
⟨100⟩
Easy
⟨111⟩
Medium
⟨110⟩
⟨100⟩
Hard
Ni
M(emu/cm³)
0
100
200
300
400
500
0
100
200
300
400
500
600
H(Oe)

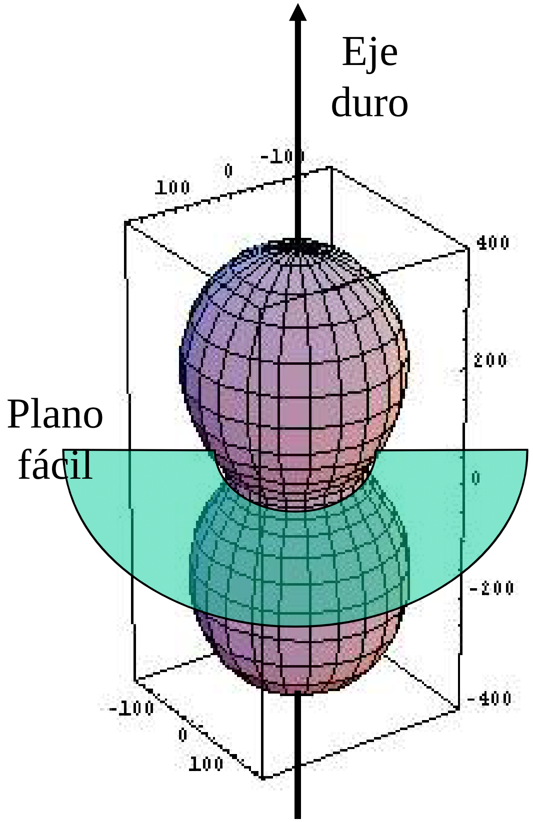
Dependencia de la anisotropía con la temperatura en la Magnetita

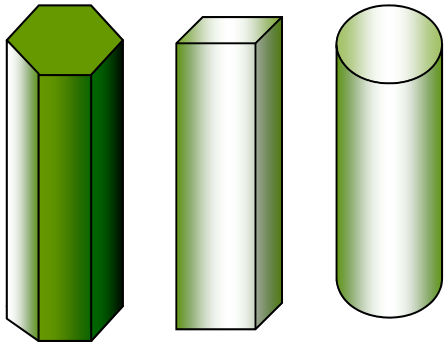


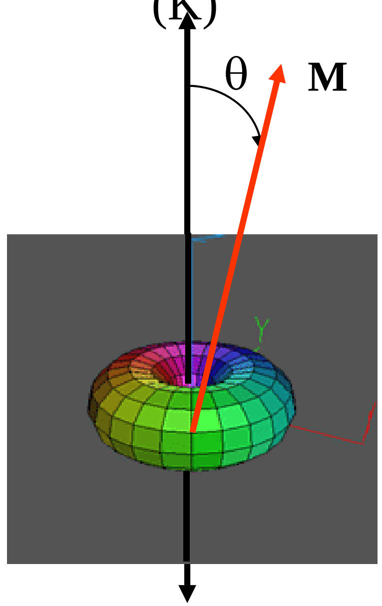


Ojo!! colores invertidos

Sistemas hexagonal y tetragonal







$$e_K = K_1 \cos^2\theta + K_2 \cos^4\theta$$

usando $\cos^2\theta = 1 - \sin^2\theta$

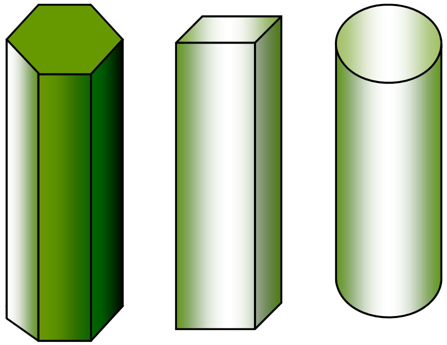
$$e_K = (K_1 + K_2) - (K_1 + 2K_2)\sin^2\theta + K_2 \sin^4\theta$$

llamando *y corriendo el cero de la energía*

$$K_1' = -(K_1 + 2K_2)$$

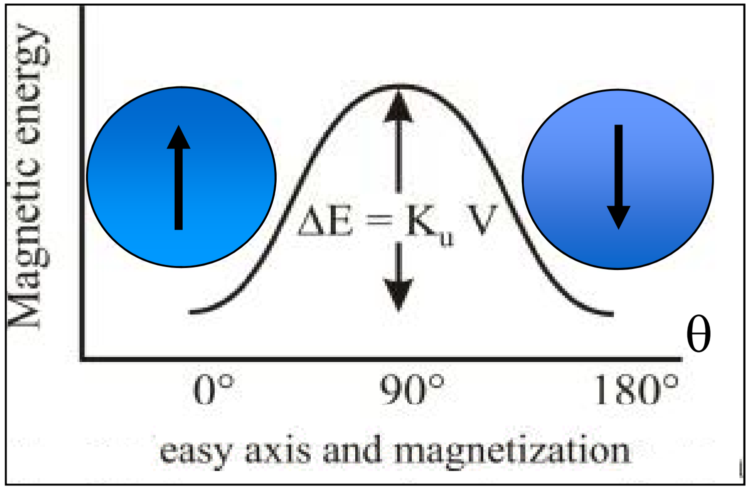
$$K_2' = K_2$$

$$e_K = K_1' \sin^2\theta + K_2' \sin^4\theta$$



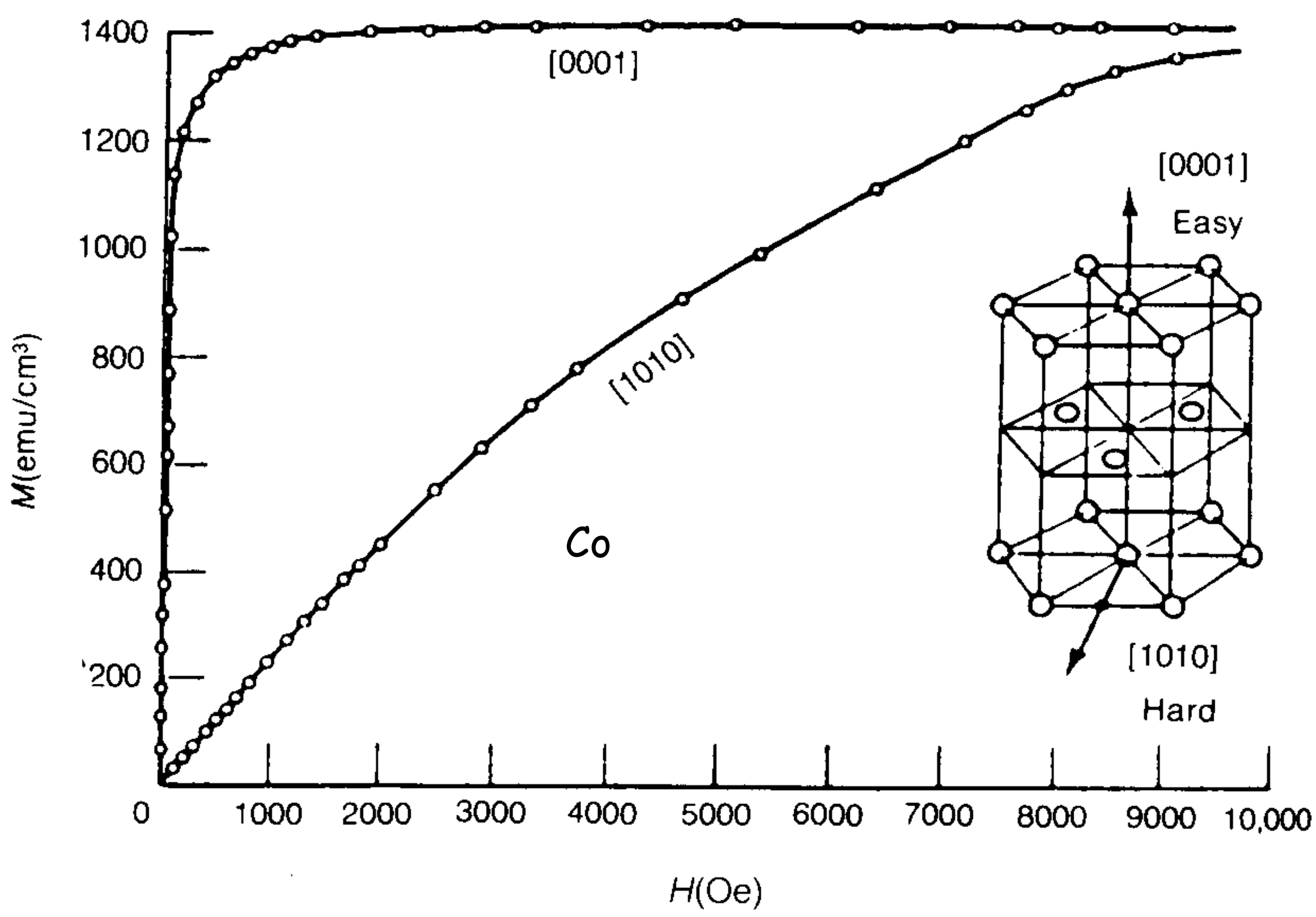
$$e_K = K_1 \sin^2\theta + K_2 \sin^4\theta \xrightarrow[\textit{caso uniaxial}]{\textit{descripción simple}} e_K = K \sin^2\theta$$

$$\boxed{E_K = KV \sin^2 \theta}$$



$$e_K = K_1 \sin^2 \theta + K_2 \sin^4 \theta$$

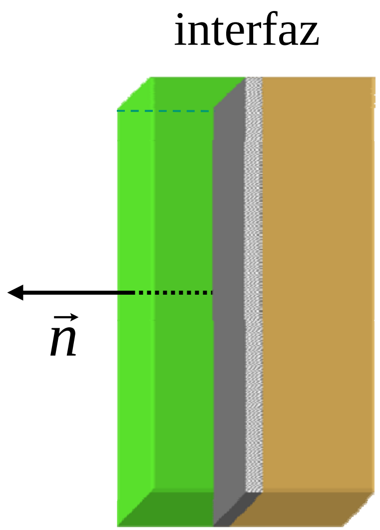
Material	K_1 (10^5 J/m^3)	K_1 (10^5 J/m^3)	Eje fácil
Co	4.1	1.0	hexagonal
$SmCo_5$	1100	-	hexagonal


1400
1200
1000
800
600
400
200
0
1000
2000
3000
4000
5000
6000
7000
8000
9000
10,000
M(emu/cm³)
H(Oe)
[0001]
[1010]
Co
[0001]
Easy
[1010]
Hard

Anisotropía de Intercambio

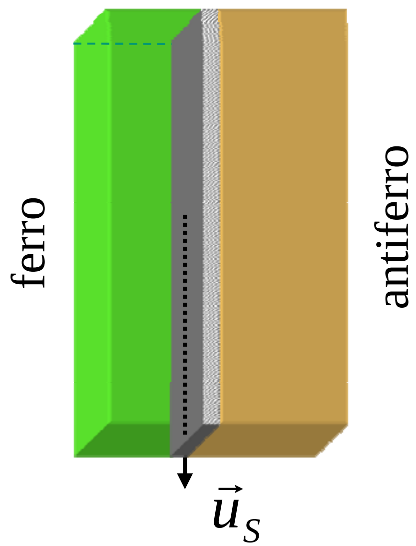
superficies e interfaces

Anisotropía de interfaz



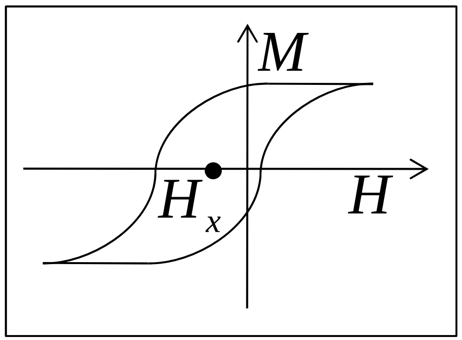
$$e_K = K_S\left[1-(\vec{m}\cdot\vec{n})^2\right] \begin{cases} K_S > 0 \Rightarrow \vec{m} \,//\, \mathrm{sup} \\ K_S < 0 \Rightarrow \vec{m} \perp \mathrm{sup} \end{cases}$$

Anisotropía de intercambio*



$$e_K = K_S \vec{m}\cdot\vec{u}_S = \frac{H_x}{2}\vec{m}\cdot\vec{u}_S$$

$$e_K = \frac{H_x}{2} m\cos\varphi$$

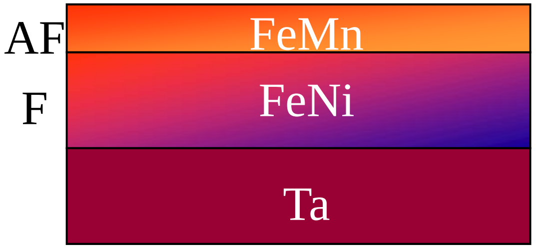


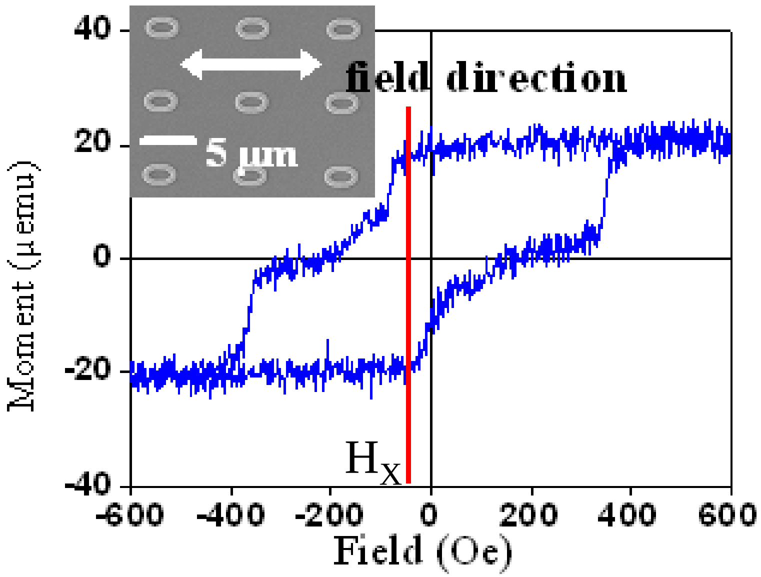
Exchange bias field

*también llamada unidireccional

Observación del exchange bias

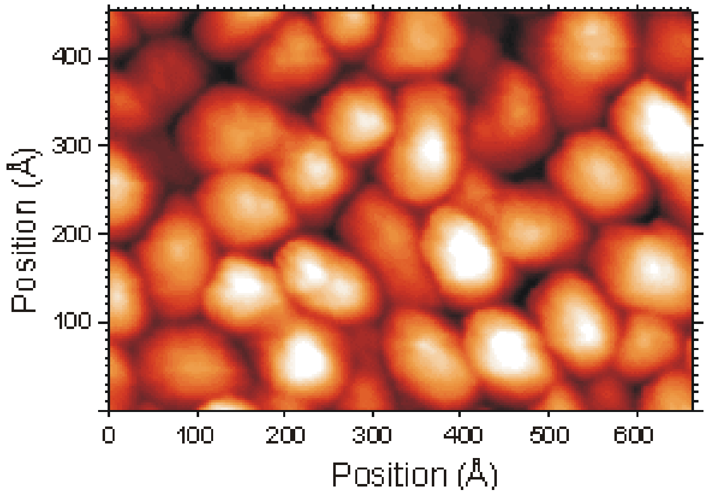
$$e_K = K_S \vec{m} \cdot \vec{u}_S = \frac{H_x}{2} \vec{m} \cdot \vec{u}_S$$

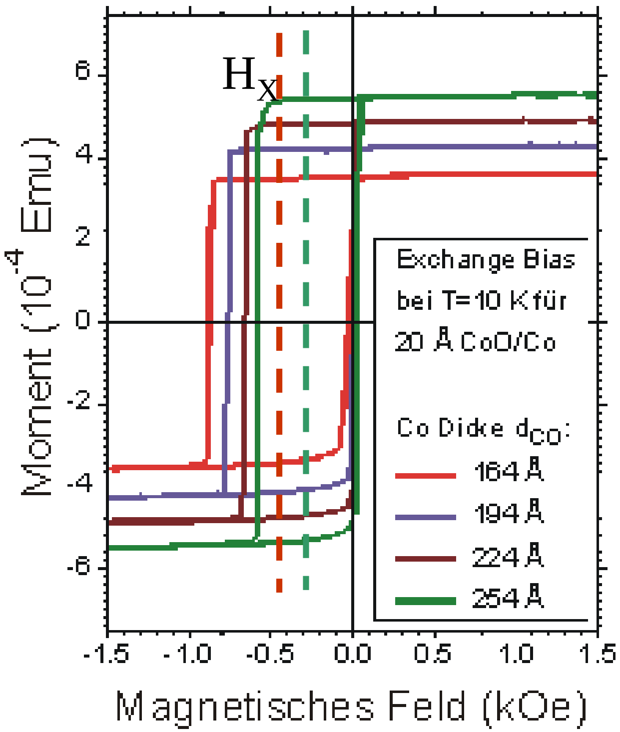


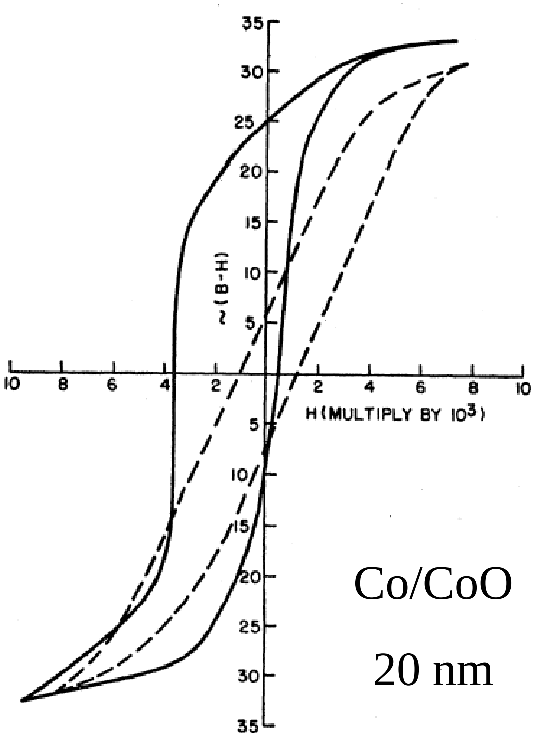


Ta 20nm / NiFe 20nm/ FeMn 10 nm film

CoO/Co



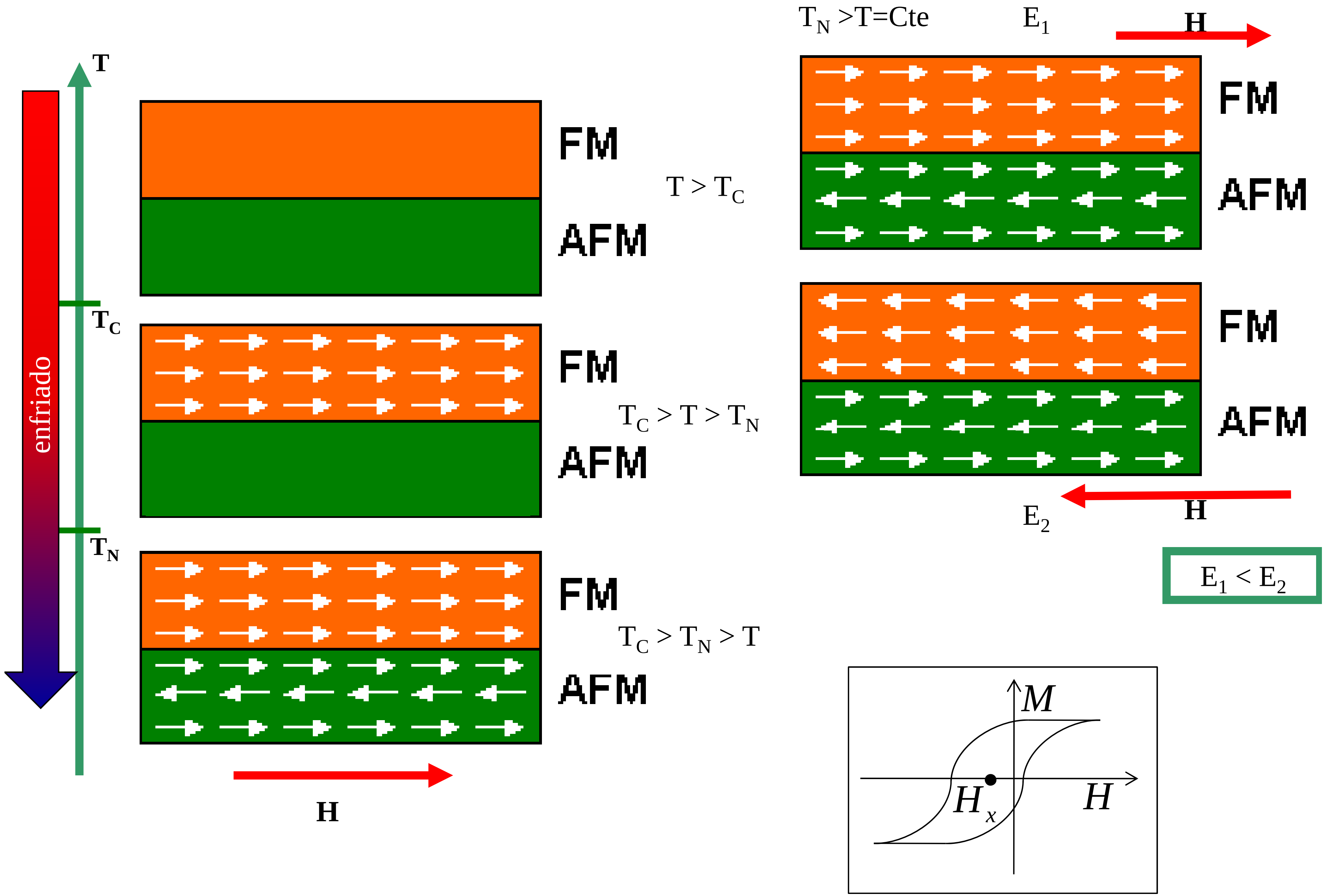




Letters to the Editor

New Magnetic Anisotropy

W. H. MEIKLEJOHN AND C. P. BEAN
General Electric Research Laboratory, Schenectady, New York
(Received March 7, 1956)
PHYSICAL REVIEW
VOLUME 102, NUMBER 5
JUNE 1, 1956



T
enfriado
T_C
T_N
FM
AFM
$T > T_C$
FM
AFM
$T_C > T > T_N$
FM
AFM
$T_C > T_N > T$
H
$T_N > T = Cte$
E_1
H
FM
AFM
FM
AFM
E_2
H
$E_1 < E_2$
M
H_x
H

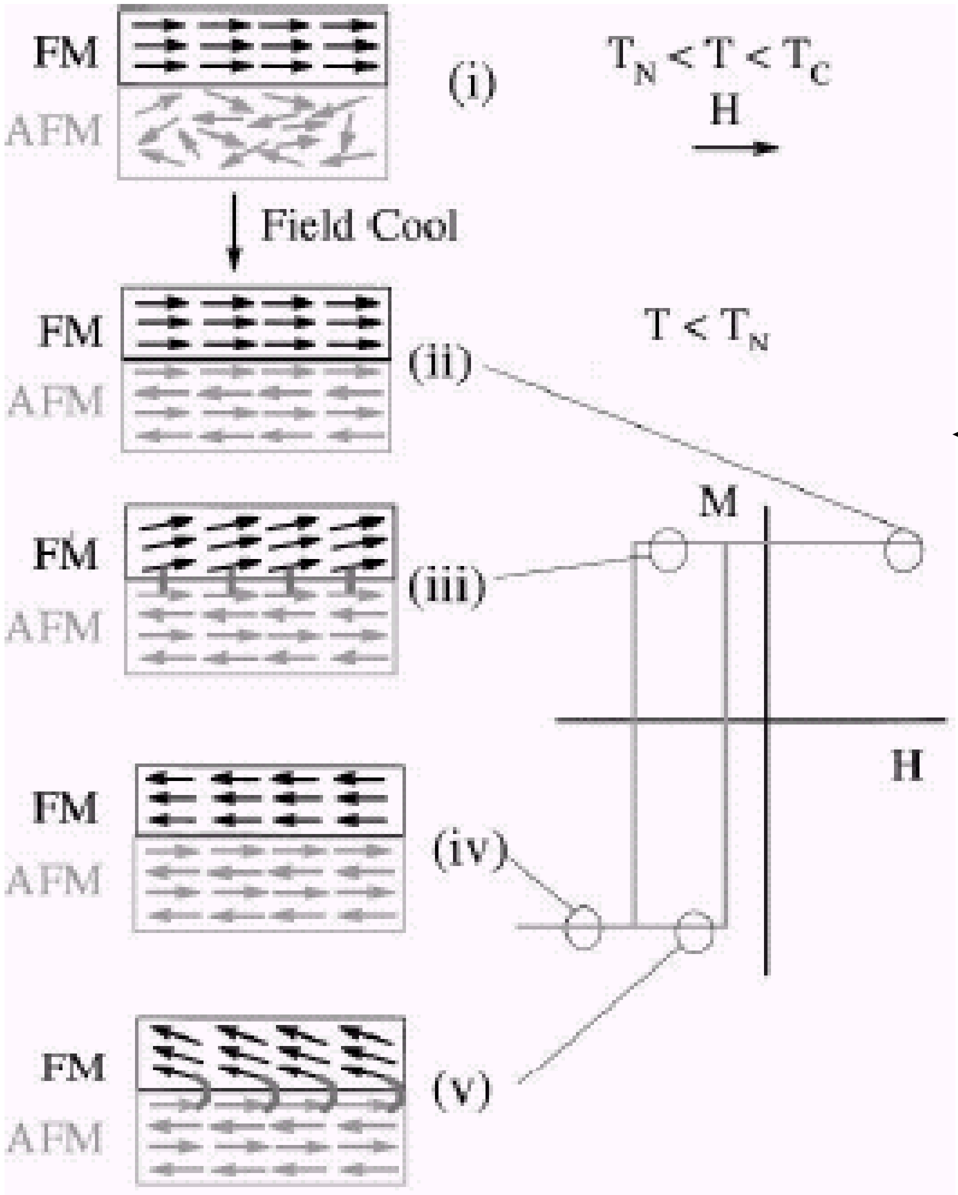
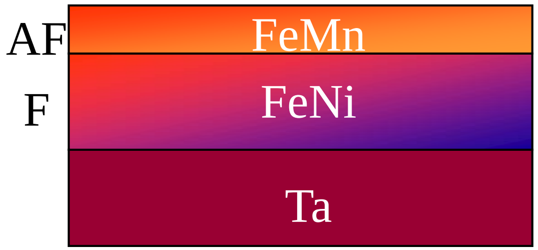
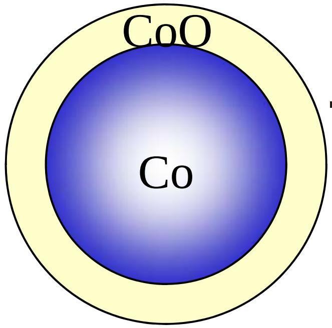


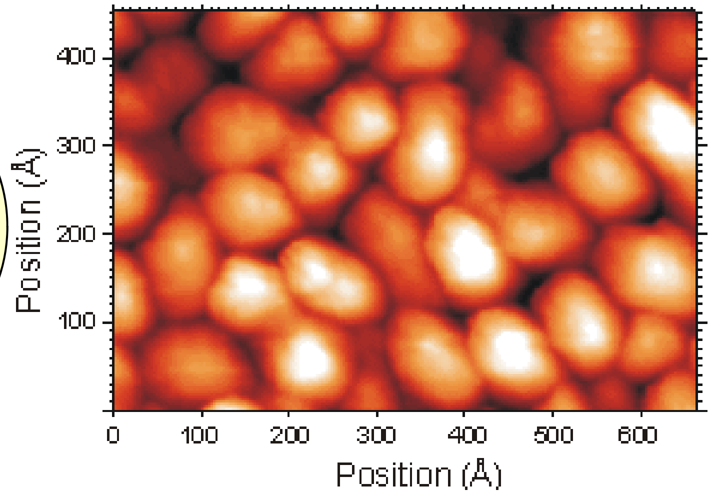
Diagrama esquemático de una configuración de spin FM /AFM.

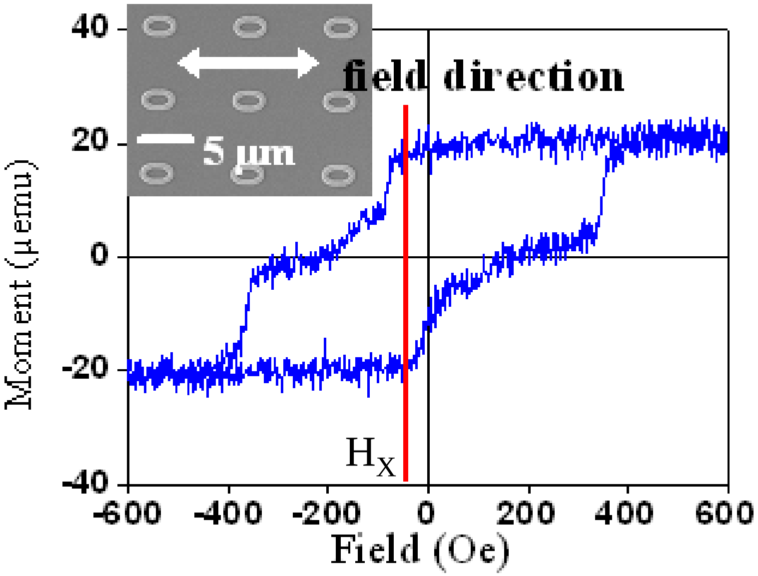
Observación de la anisotropía de intercambio (Exchange bias)



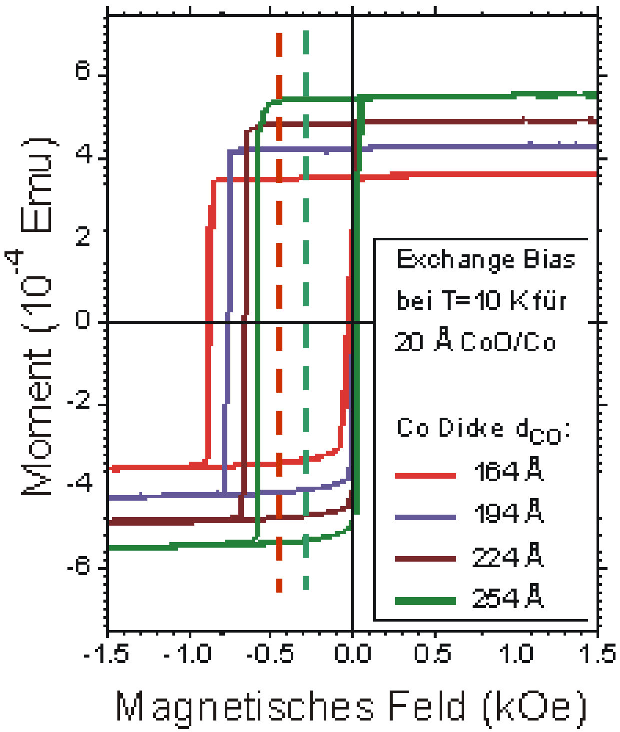


CoO/Co



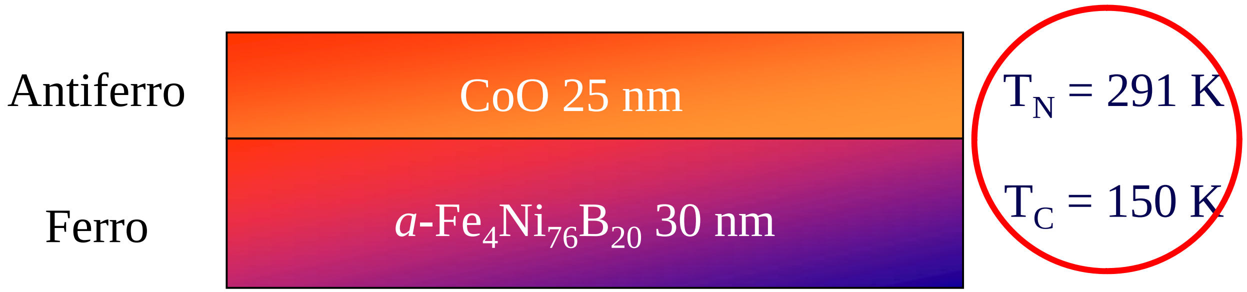


Ta 20nm / NiFe 20nm/ FeMn 10 nm film



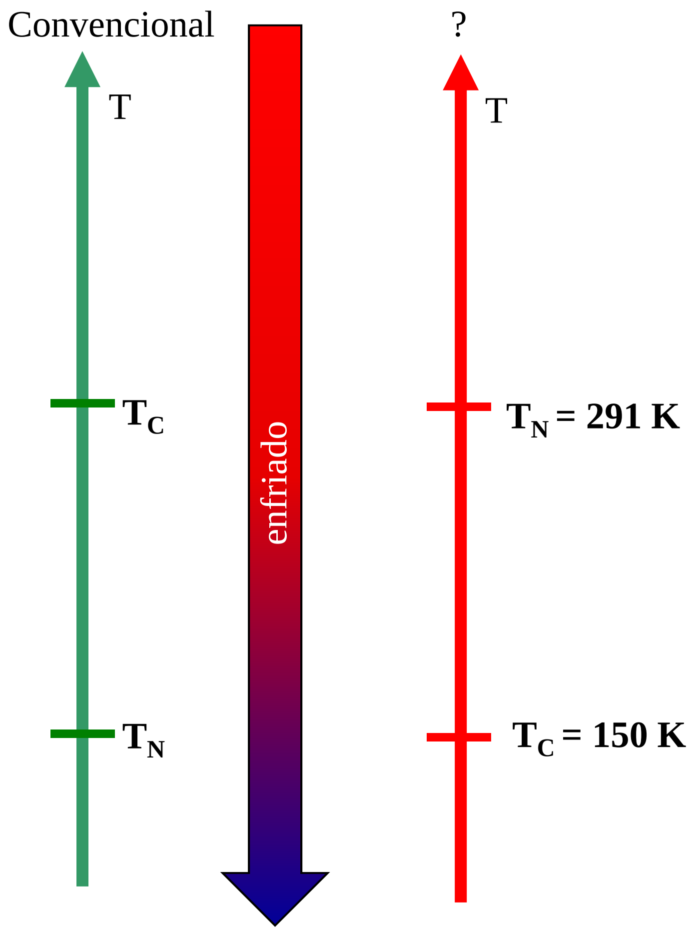
Exchange Coupling in the Paramagnetic State

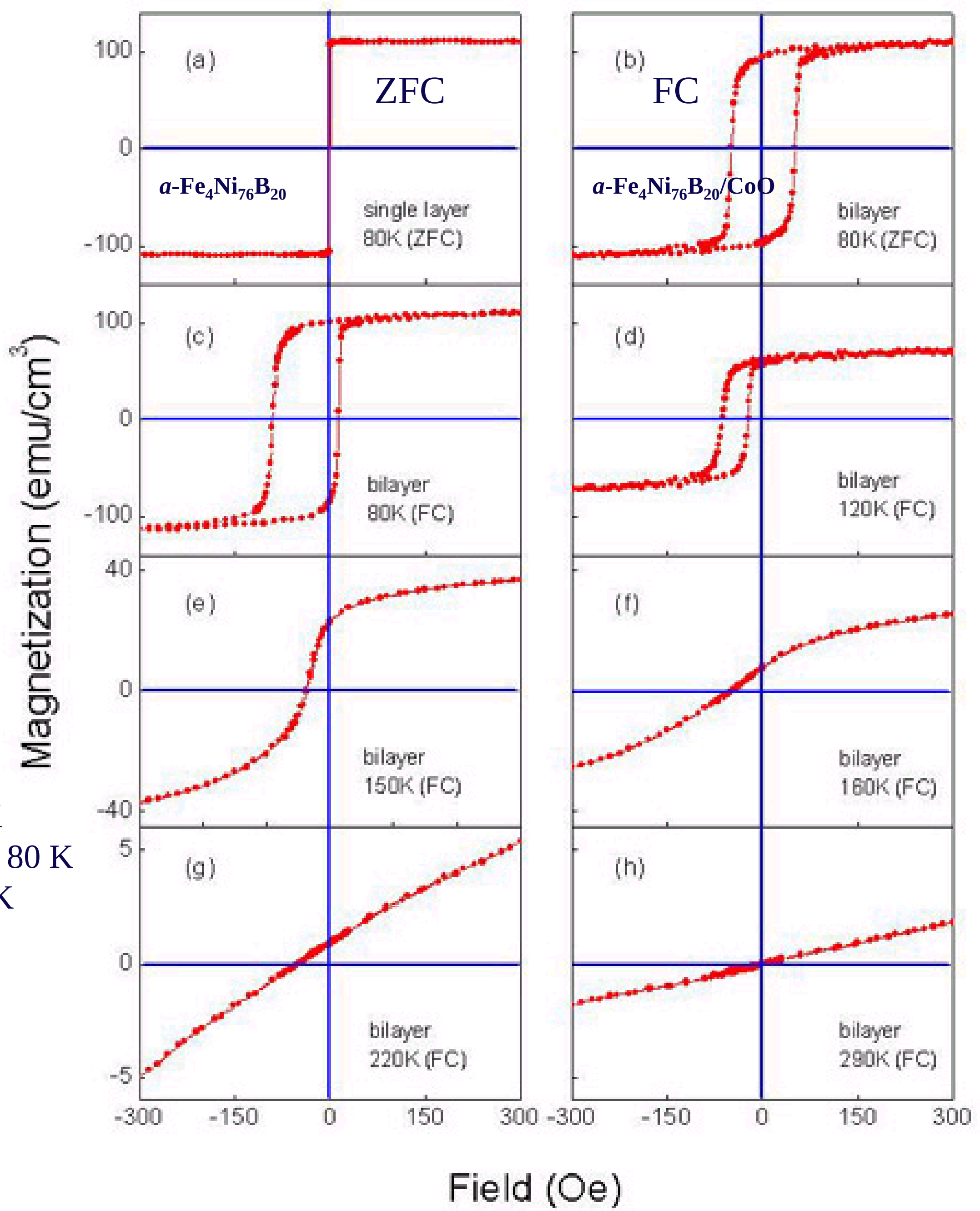
J. W. Cai, Kai Liu, and C. L. Chien, *The Johns Hopkins University, Baltimore, MD 21218*



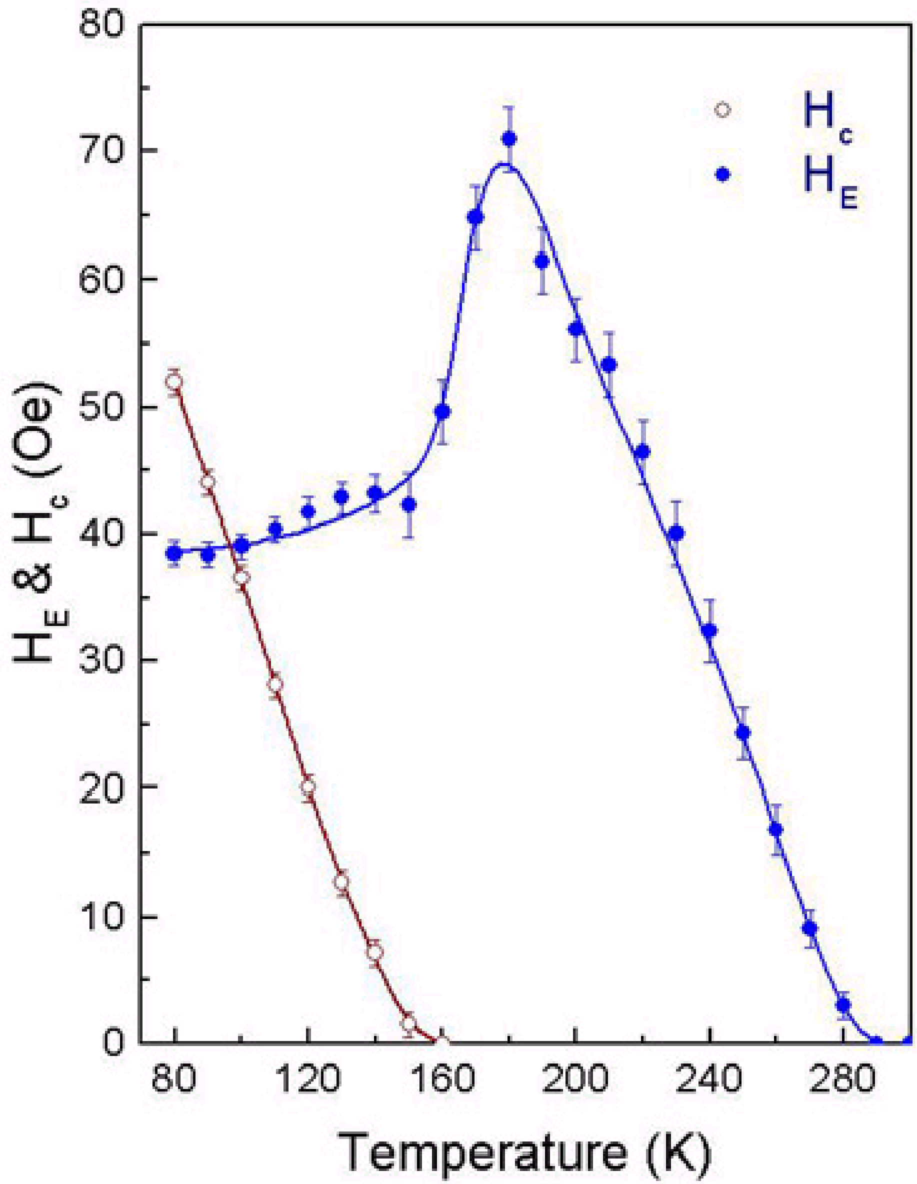
Reference

J. W. Cai, Kai Liu, and C. L. Chien, Phys. Rev. B 60, 72 (1999).





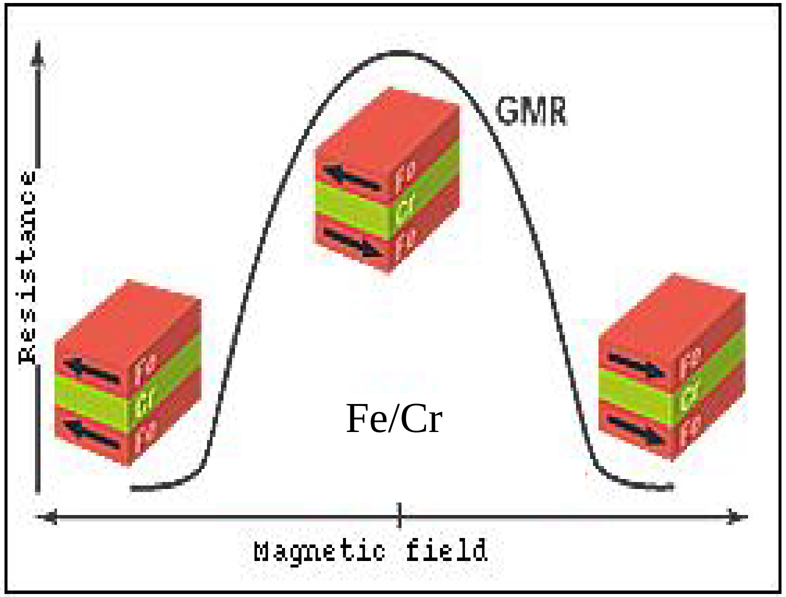
(a), (b) ZFC hasta 80 K
(c) FC en 10 kOe hasta 80 K
(d) Llevado hasta 120 K
(e) 150 K
(f) 160 K
(g) 220 K (g)
(h) 290 K (h)

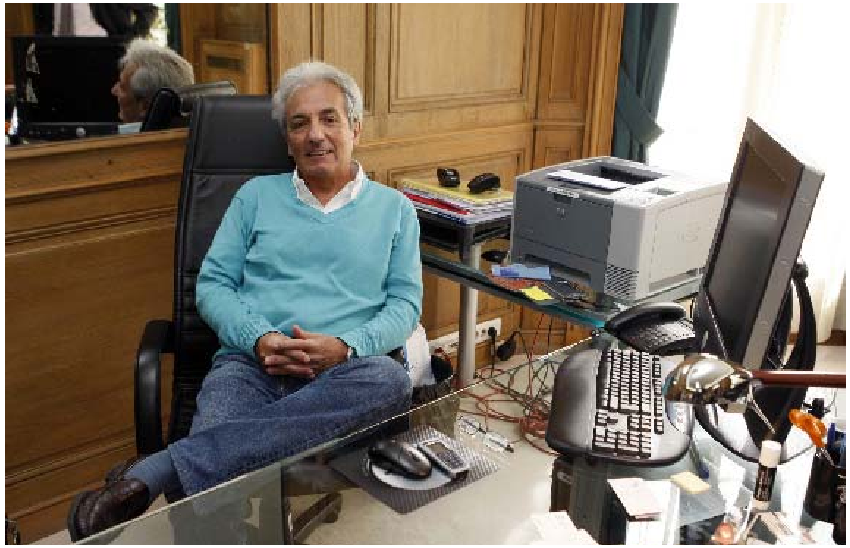


Temperature dependence of exchange field H_E and coercivity H_C of *a*-$Fe_4Ni_{76}B_{20}$(30 nm)/CoO (25 nm) after field cooling in 10 kOe to 80 K.

Anisotropía de intercambio – válvula de spin

Magnetoresistencia gigante





Albert Fert, Nobel Prize in Physics 2007

Giant Magnetoresistance of (001) Fe/(001) Cr Magnetic Superlattices

M. N. Baibich,[a] J. M. Broto, A. Fert, F. Nguyen Van Dau, and F. Petroff
Laboratoire de Physique des Solides, Université Paris-Sud, F-91405 Orsay, France

P. Eitenne, G. Creuzet, A. Friederich, and J. Chazelas
Laboratoire Central de Recherches, Thomson CSF, B.P. 10, F-91401 Orsay, France

(Received 24 August 1988)

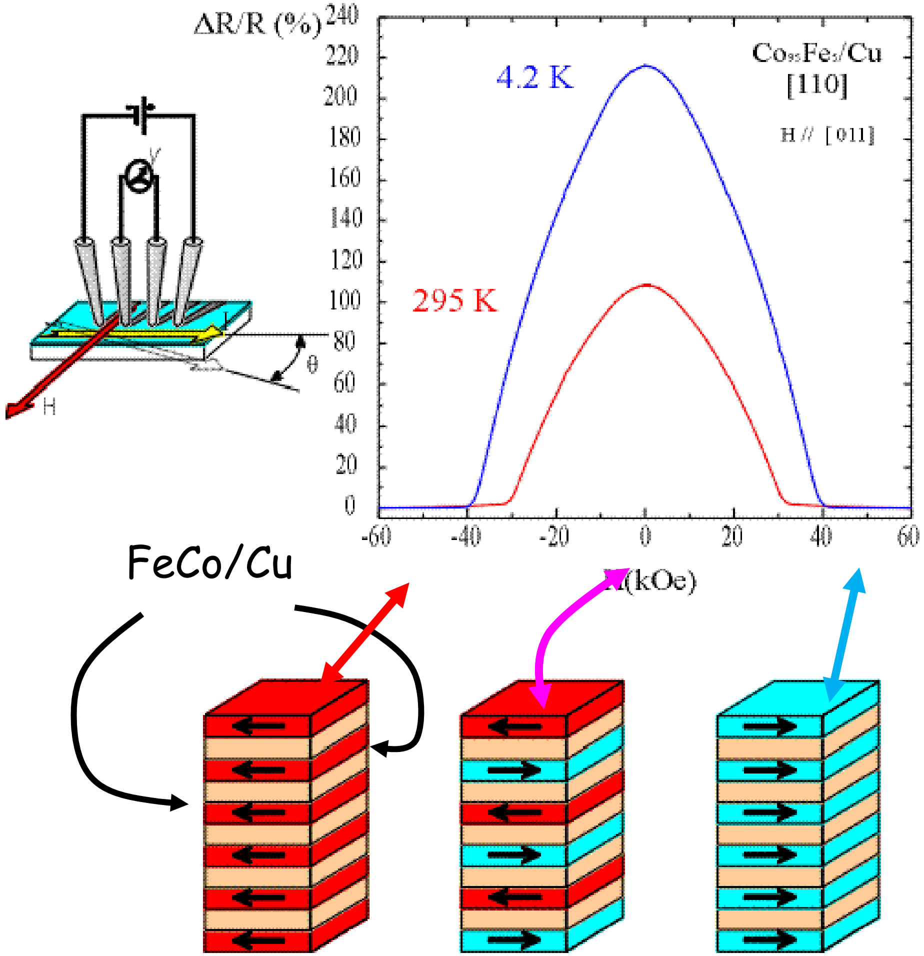
PHYSICAL REVIEW LETTERS

VOLUME 61, NUMBER 21

21 NOVEMBER 1988

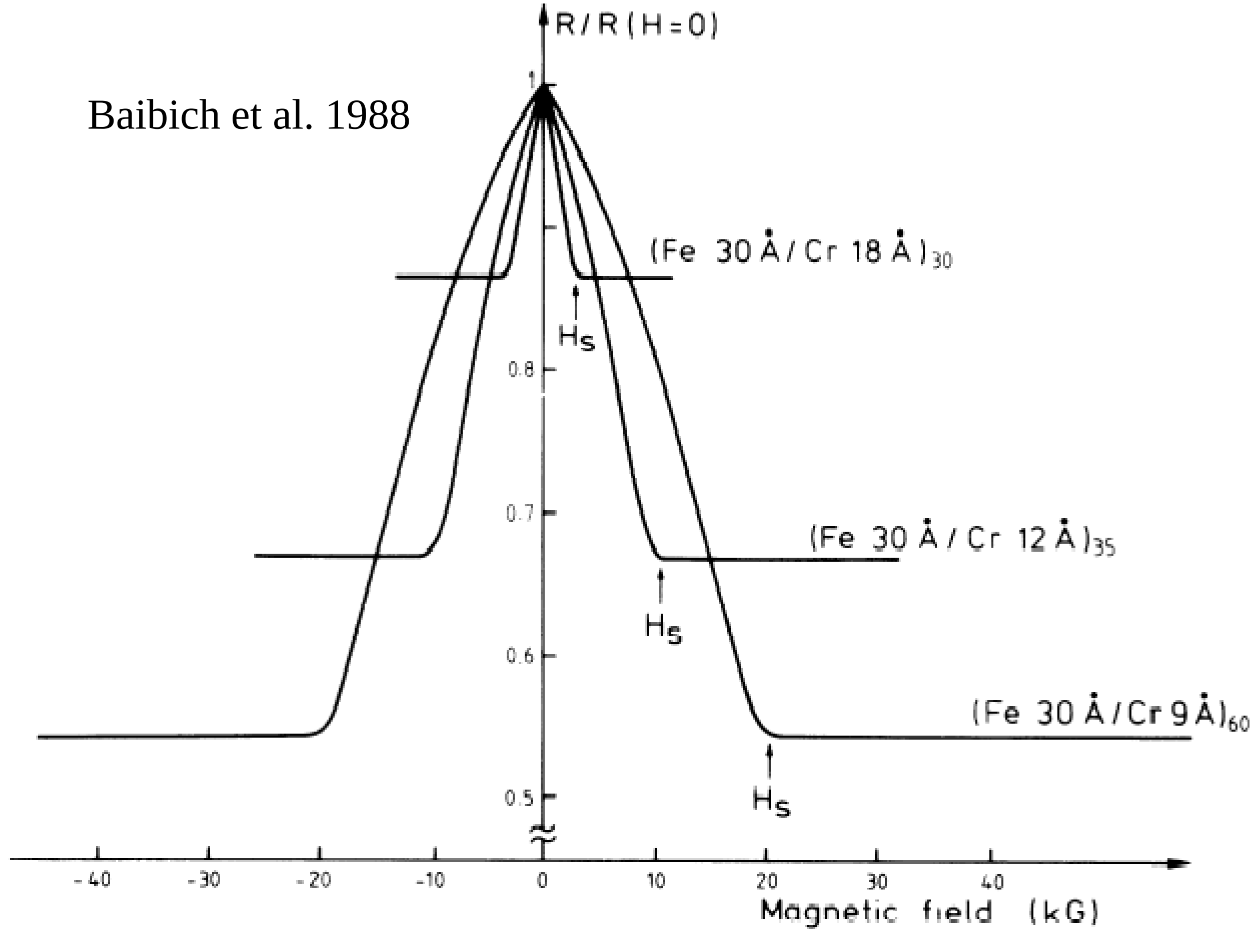
Peter Grünberg, Nobel Prize in Physics 2007

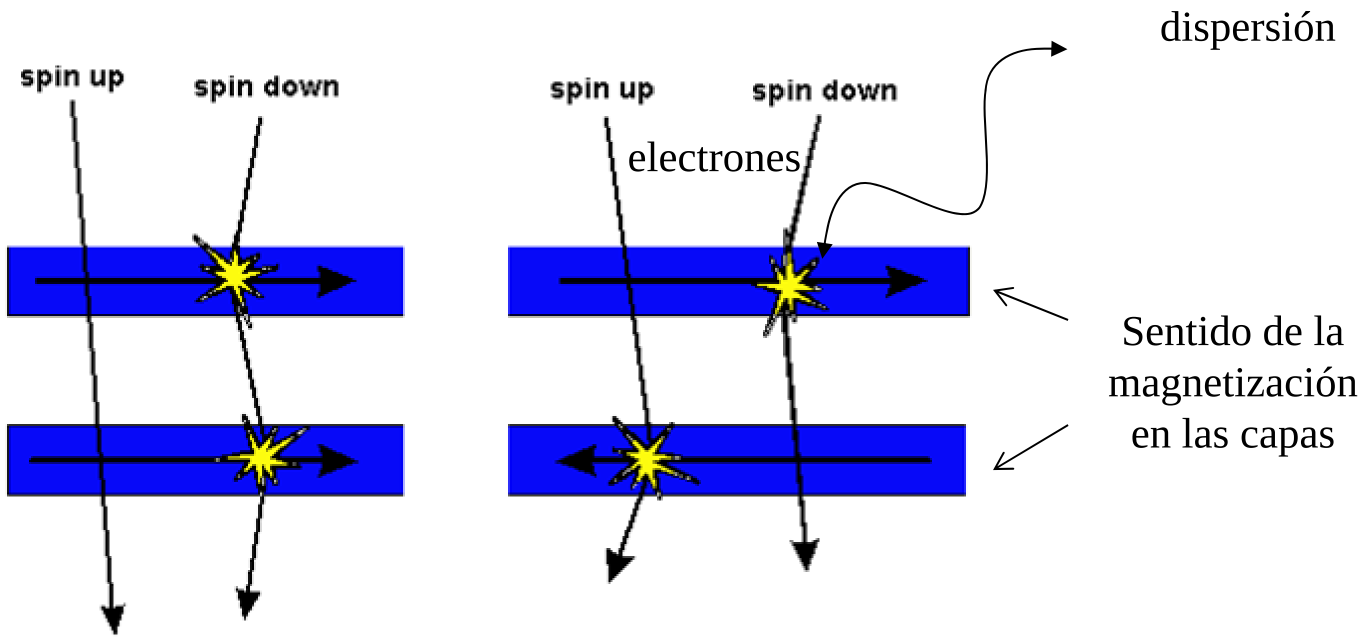
Giant Magnetoresistance



Resultados experimentales

Baibich et al. 1988





Acoplamiento ferro

Baja resistencia

Acoplamiento antiferro

Alta resistencia

Magnetoresistance - spin valve

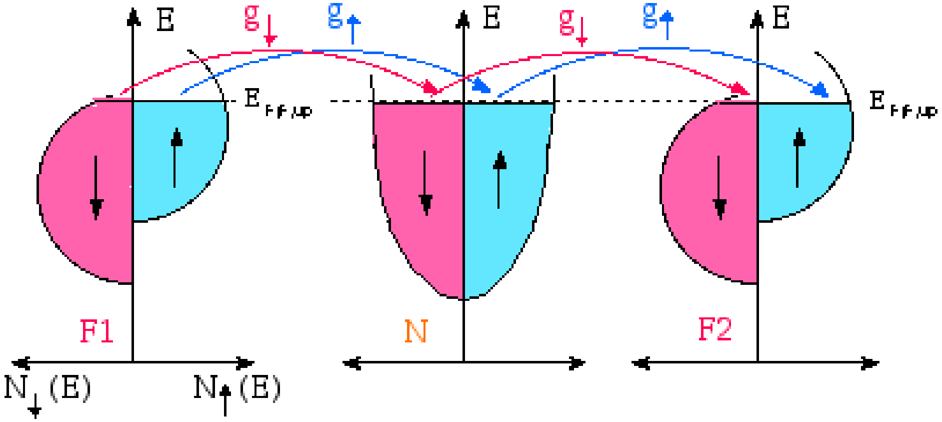
$g \propto N(E_F)$

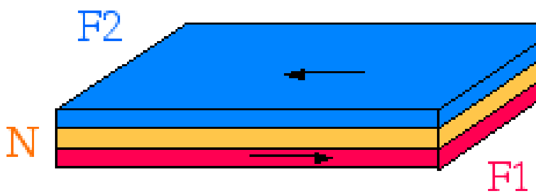
F1 and F2 parallel:

$g_{\uparrow} > g_{\downarrow}$

$g_{\uparrow}$ dominates

Low R state

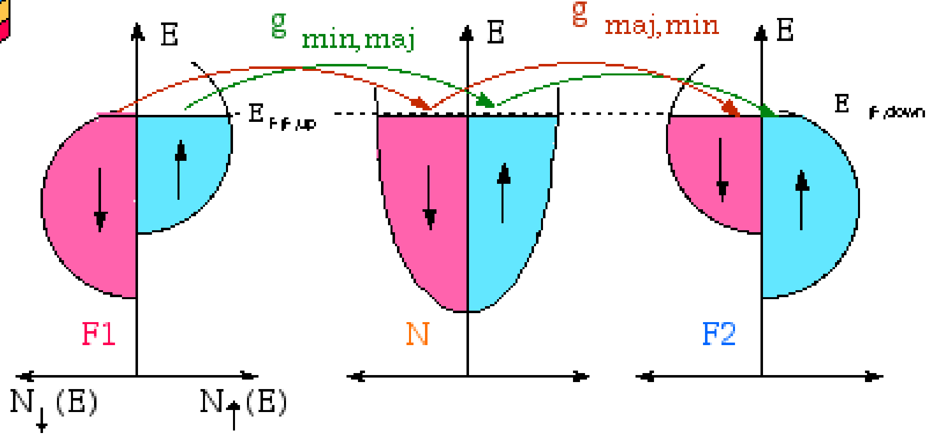




F1 and F2 antiparallel:

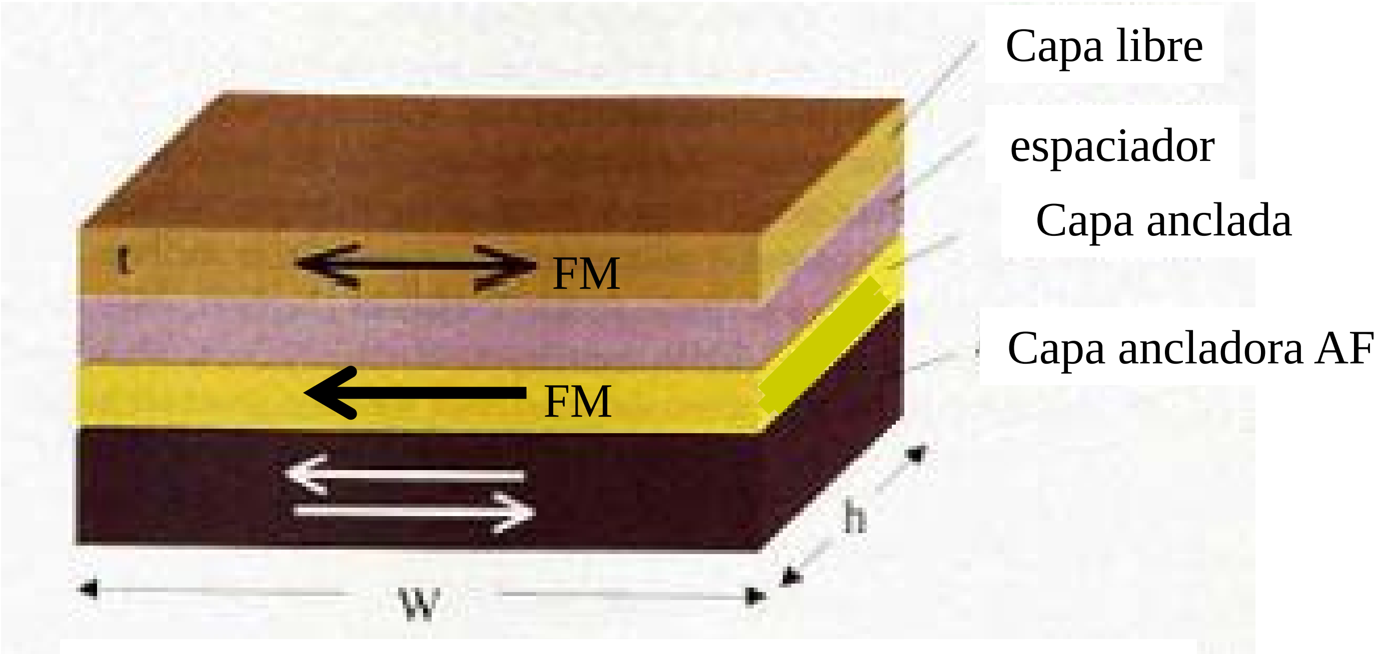
$g_{min,maj} = g_{maj,min} < g_{\uparrow}$

High R state

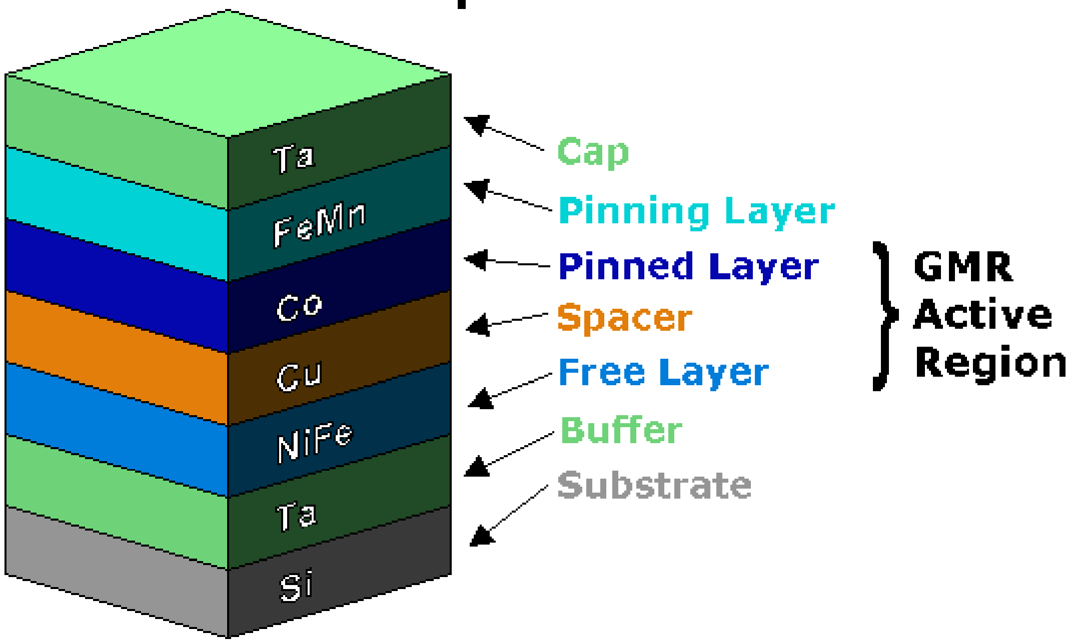


Grunberg PRB (89), Fe/Cr spin valve MR=1.5%

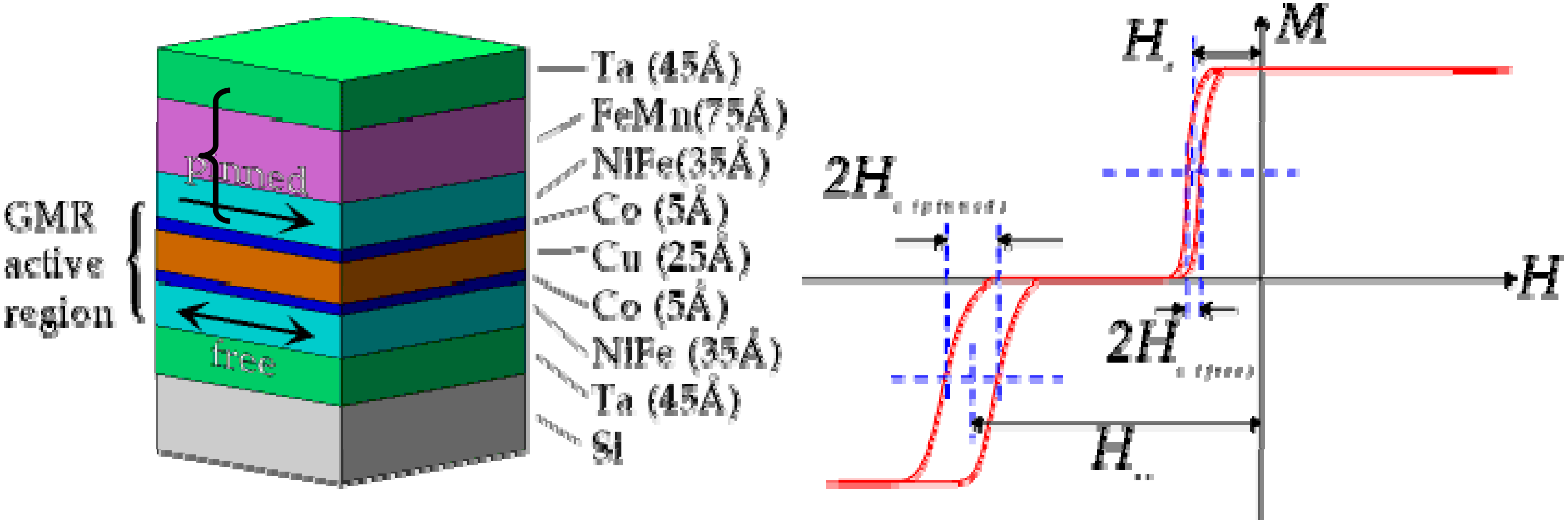
Baibich et al., PRL (88), Fe/Cr multilayer

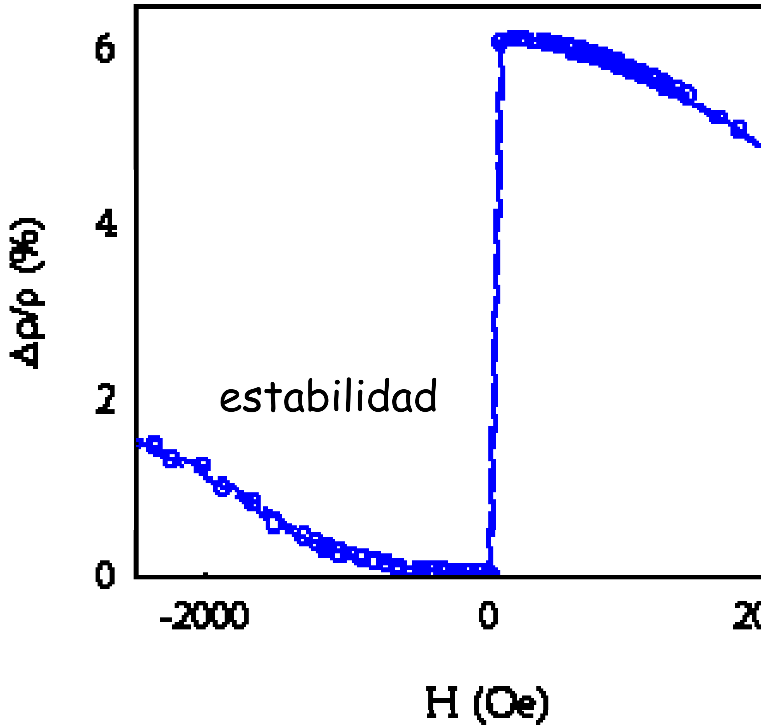


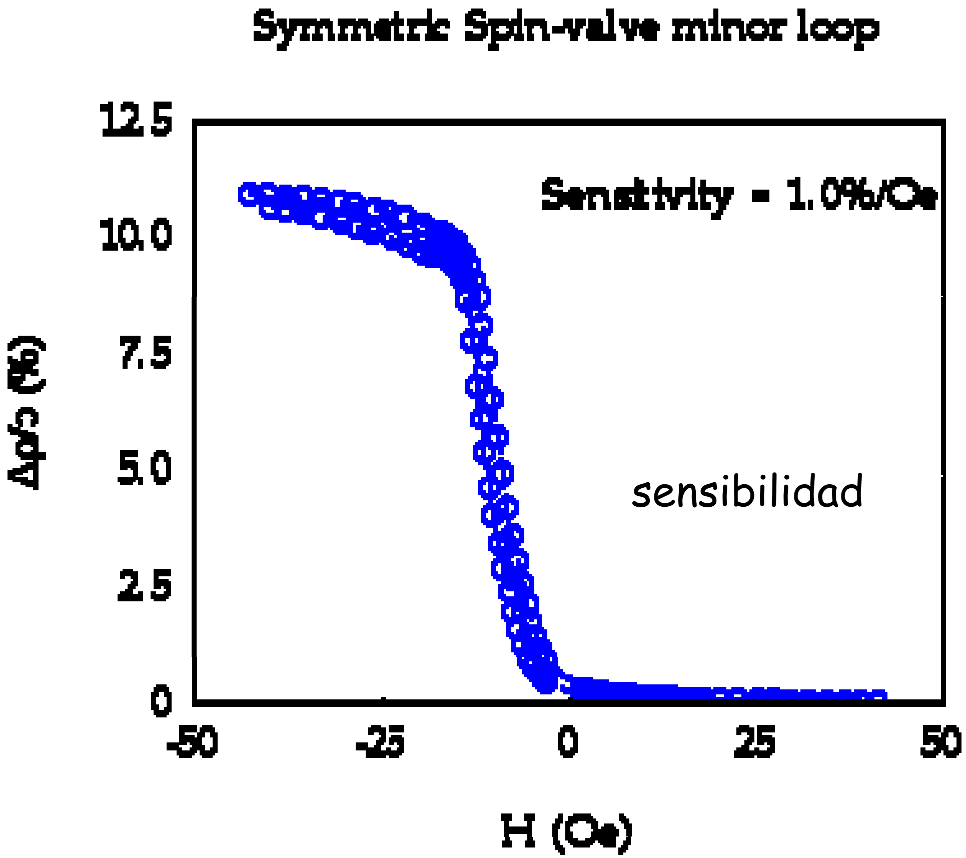
Spin Valve Structure



Válvula de spin



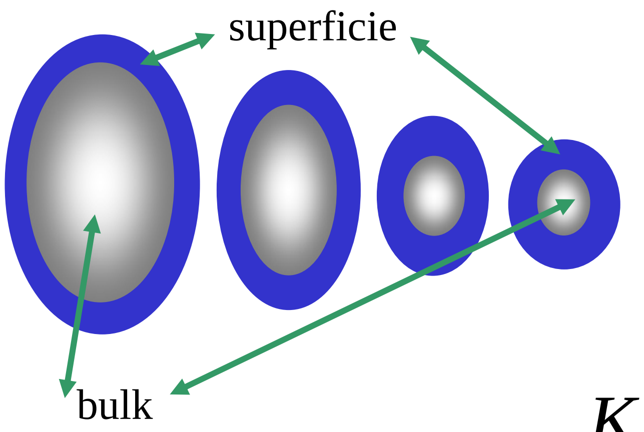

6
4
2
0
-2000
0
H (Oe)
Δρ/ρ (%)
estabilidad


Symmetric Spin-valve minor loop
12.5
10.0
7.5
5.0
2.5
0
-50
-25
0
25
50
H (Oe)
Δρ/ρ (%)
Sensitivity = 1.0%/Oe
sensibilidad

animación

Anisotropía en nanopartículas magnéticas y fluctuaciones térmicas

anisotropía de superficie en nanopartículas



$$K_{ef} = K_B + K_{V_S}^{ef}$$

K_S = 10^{-4}-10^{-3} J/m^2, energía de anisotropía por unidad de área superficial

$$K_{V_S}^{ef} \approx \frac{SK_S}{V} \approx \frac{4\pi r^2}{\frac{4}{3}\pi r^3} K_S = \frac{3}{r} K_S = \frac{6}{d} K_S$$

Partícula esférica

$$K_{ef} = K_B + \frac{6K_s}{d}$$

superficies/interfaces:

- discontinuidad composicional y configuracional
- mayor efecto anisotrópico

$$\boxed{K_{ef} = K_B + \gamma \frac{K_s}{\bar{d}}}$$

Bødker et. Al (1994)

Anisotropía de superficie - ejemplo

$$K_B\left(Co_{fcc}\right)\approx 1\times10^5\,J/m^3$$

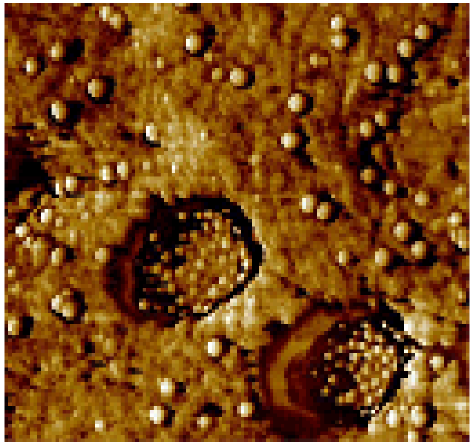


imagen MFA de nanopartículas de Co fcc en una matriz de alúmina. Las partículas son de aprox 11 nm (diámetro).

Alúmina

$$K_S\left(Co/Al_2O_3\right)\approx 3.3x10^{-4}\,J/m^2$$

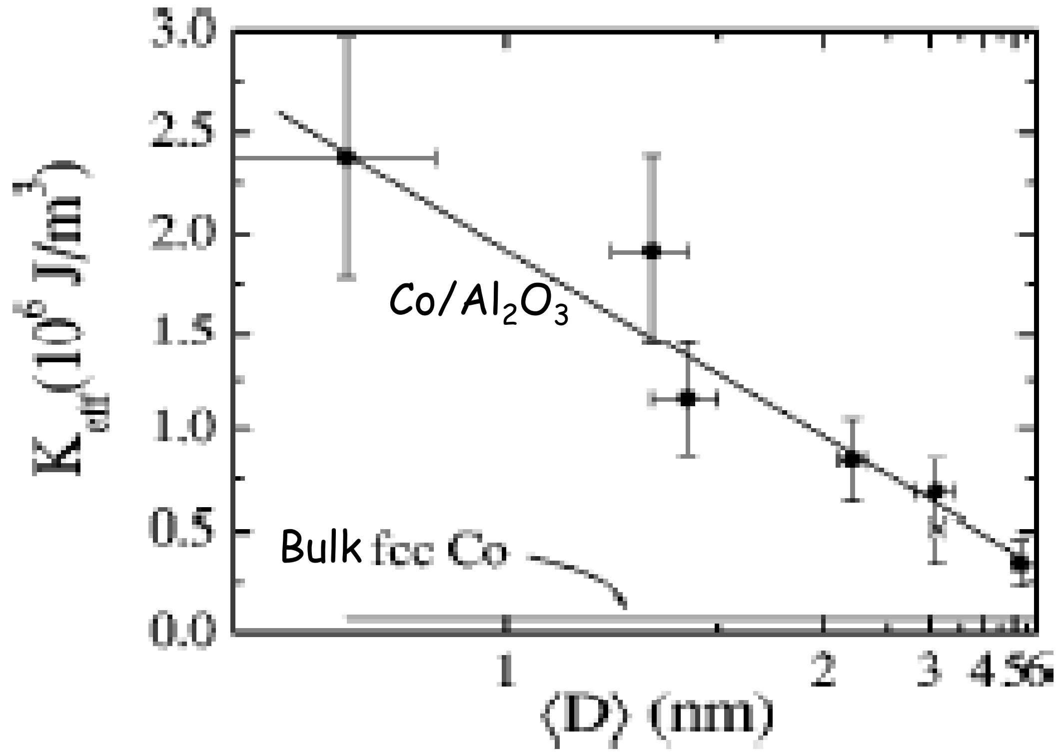
$$K_{ef}=K_B+\gamma\frac{K_s}{\bar{d}}$$

$$K_{ef}\left(Co/Al_2O_3\right)\approx\left[1\times10^5+6\frac{3.3\times10^{-4}}{11\times10^{-9}}\right]J/m^3\approx 2.8\times10^5\,J/m^3$$

Si d ~ 3 nm = 3x10^{-9}m $\Rightarrow$ $K_{ef}\left(Co/Al_2O_3\right)\approx 10^6\,J/m^3$

$$\tau=\tau_0 e^{\frac{K_{ef}V}{kT}}$$

Mayores tiempos de relajación



F. Luis, J.M. Torres, L.M. Gracía, J. Bartolomé, J. Stankiewicz, F. Petroff, F. Fettar, J. L. Maurice and A. Vaurés. Phys. Rev B, **65** (2002) 094409

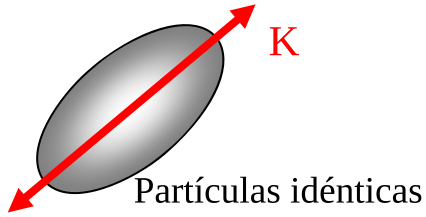
Modelo de Stoner - Wohlfarth

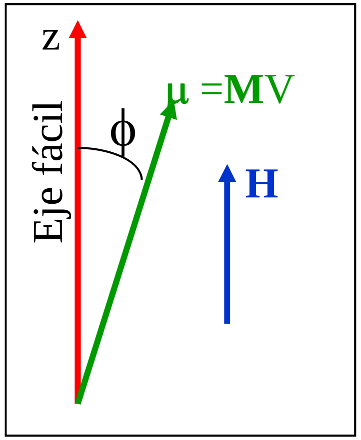
monodominio

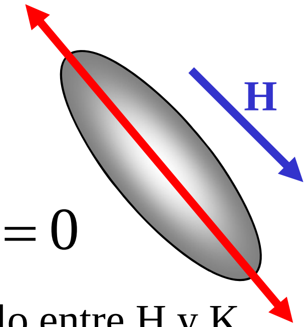
Anisotropía uniaxial

no interactuantes

Campo paralelo al eje fácil







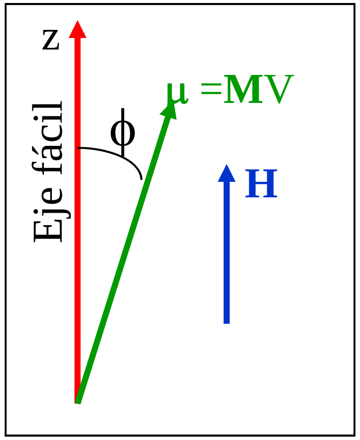
$$\theta = 0$$

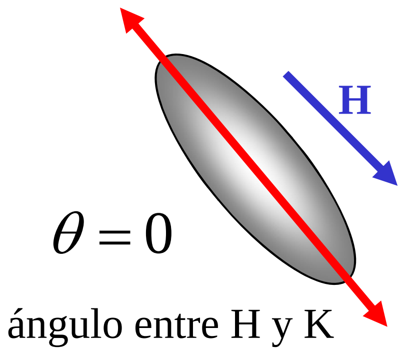
ángulo entre H y K

$$E_K = e_K V = KV \sin^2 \phi$$

$$E_H = -\vec{\mu} \cdot \vec{B} = -\mu_0 \vec{\mu} \cdot \vec{H} = -\mu_0 V M_z H = -\mu_0 V M_S H \cos\phi$$

$$E = E_K + E_H = KV \sin^2 \phi - \mu_0 V M_S H \cos\phi$$



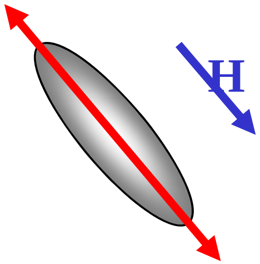


$$E = E_K + E_H = KV\sin^2\phi - \mu_0 V M_S H\cos\phi$$

llamamos Campo de anisotropía

$$\boxed{H_K = \frac{2K}{\mu_0 M_S}} \quad \boxed{h = \frac{H}{H_K} = \frac{\mu_0 M_S H}{2K}}$$

$$\boxed{E = KV\left(\sin^2\phi - 2h\cos\phi\right)}$$



$$E = KV\left(\sin^2\phi - 2h\cos\phi\right)$$

extremo
$$\frac{\partial E}{\partial \phi} = 0 \Rightarrow \sin\phi\left(\cos\phi + h\right) = 0 \Rightarrow \begin{cases} \phi = \begin{cases} 0 \\ \pi \end{cases} \\ \cos\phi = -h \end{cases}$$

mínimo
$$\frac{\partial^2 E}{\partial \phi^2} = 2KV\left[\cos\phi\left(\cos\phi + h\right) - \sin^2\phi\right] > 0$$

Valores de la función

$$E(\phi = 0) = -2hKV$$

$$E(\phi = \pi) = 2hKV$$

$$E(\cos\phi = -h) = KV(1 + h^2)$$

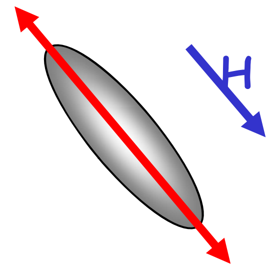
Condición de mínimo

$$\frac{\partial^2 E}{\partial \phi^2}(\phi = 0) = 2KV(1 + h) > 0$$

$$\frac{\partial^2 E}{\partial \phi^2}(\phi = \pi) = 2KV(1 - h) > 0$$

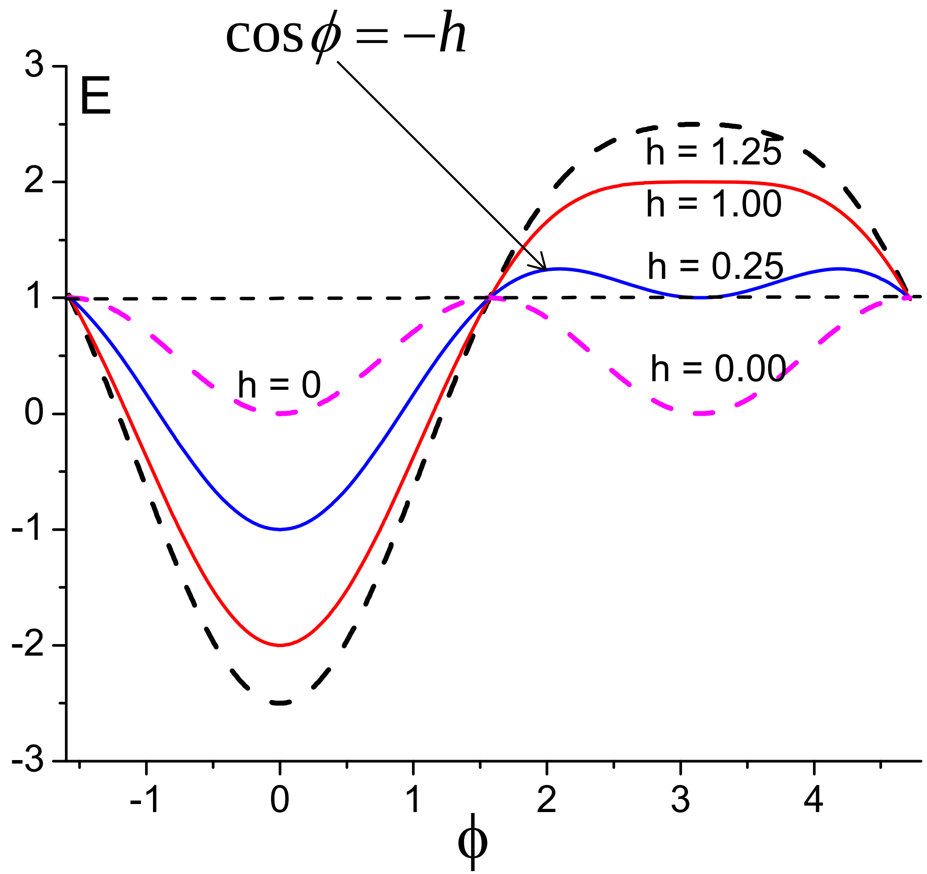
$$\frac{\partial^2 E}{\partial \phi^2}(\cos\phi = -h) = 2KV(h^2 - 1) > 0$$

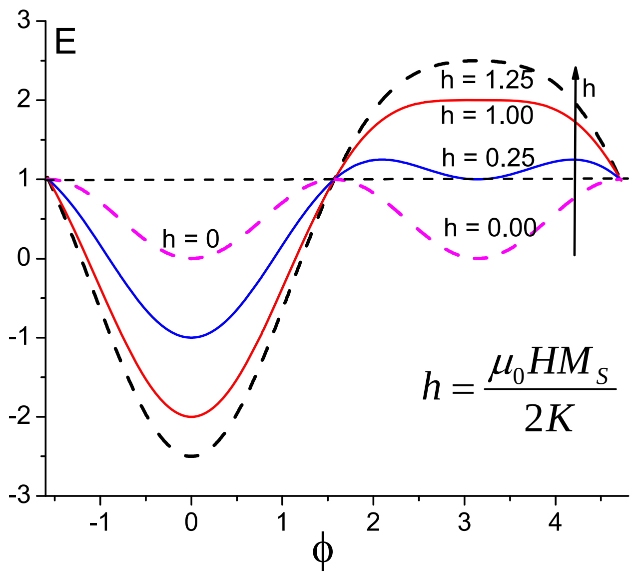
$$\textit{siempre que } h < 1 \Rightarrow \textit{máximo}$$

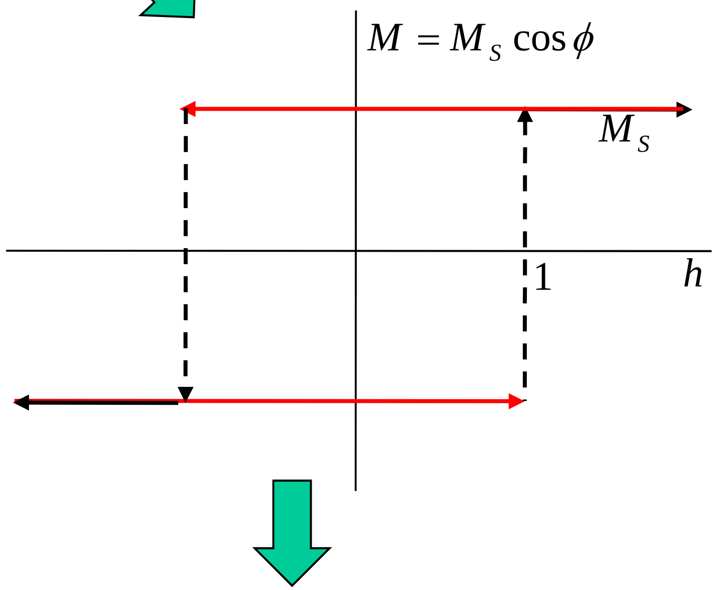


$$E = KV\left(\sin^2\phi - 2h\cos\phi\right)$$

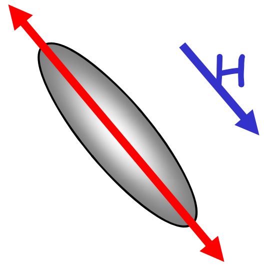
$$h = \frac{H}{H_K} \quad H_K = \frac{2K}{\mu_0 M_S}$$

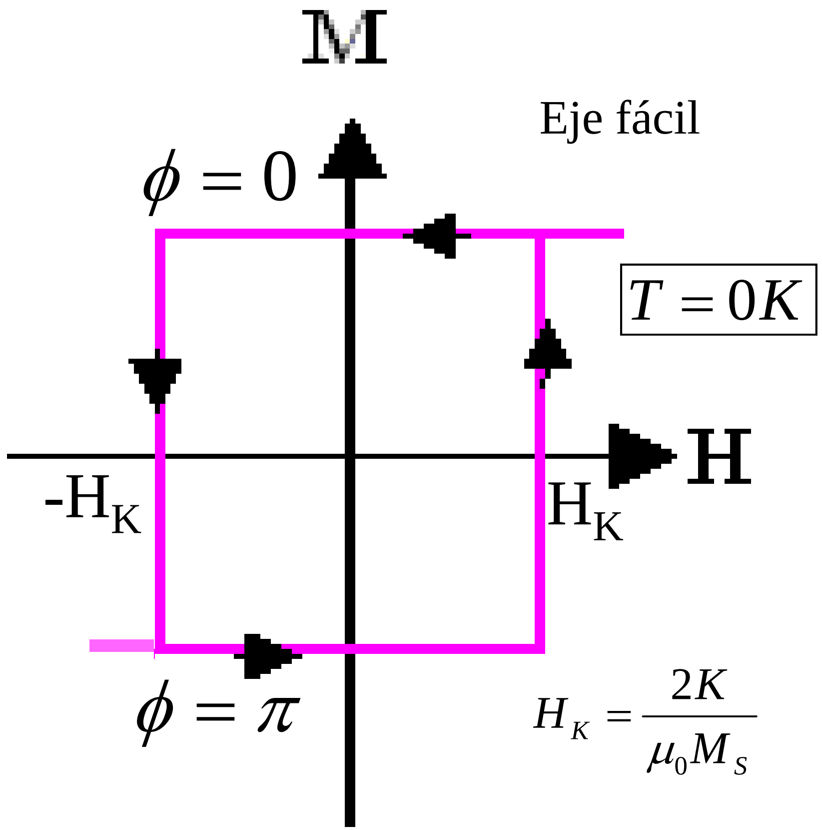



E
h = 1.25
h = 1.00
h = 0.25
h = 0
h = 0.00
h
$h = \frac{\mu_0 H M_S}{2K}$
ϕ


$M = M_S \cos\phi$
M_S
1
h

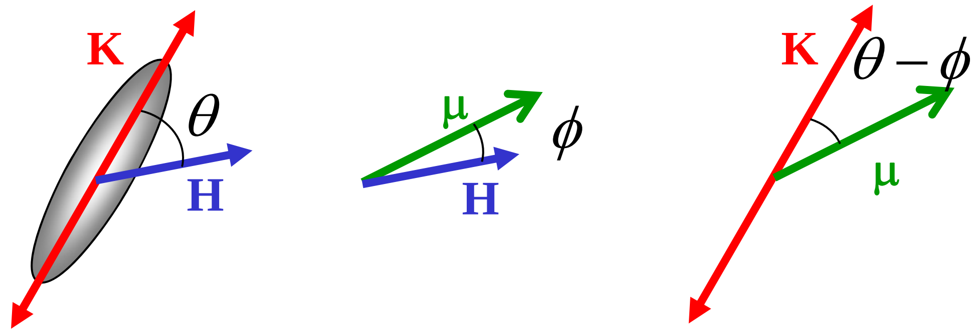
$$M_z = M_s \cos\theta$$



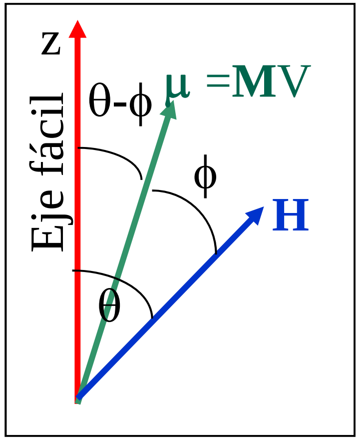


Campo en dirección arbitraria

$$\theta \neq 0$$



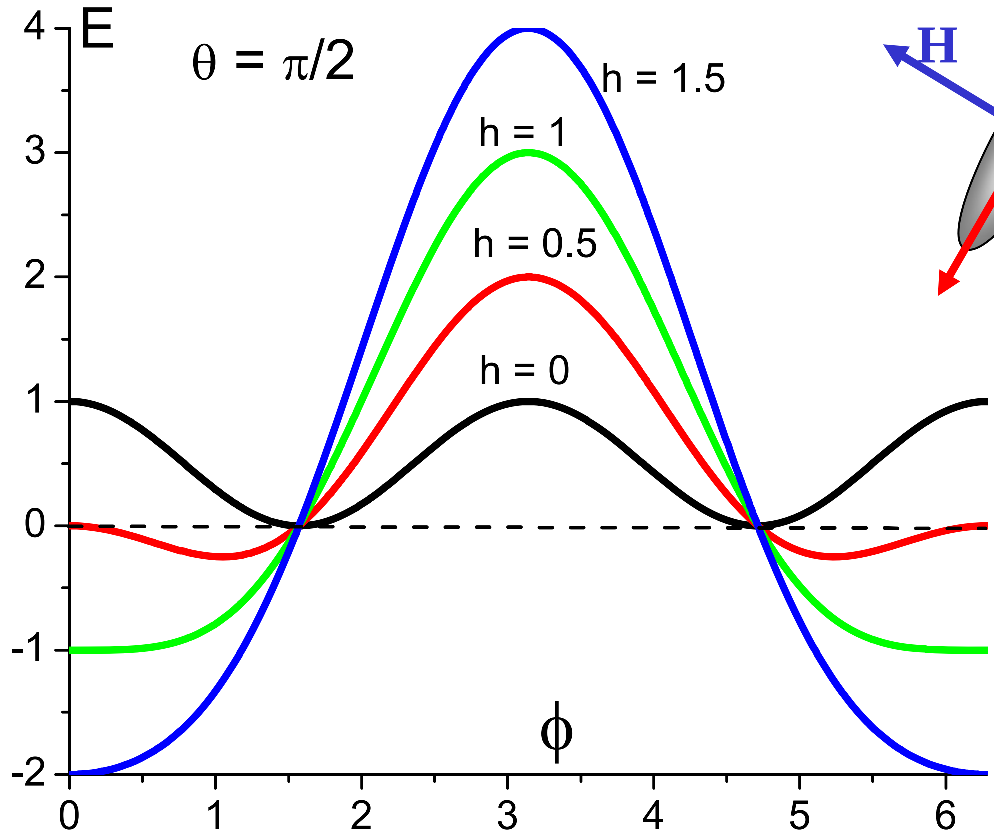
$$E = E_K + E_H = KV\left[\sin^2(\phi - \theta) - 2h\cos\phi\right]$$

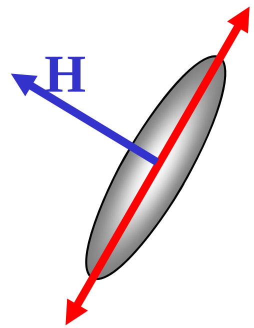


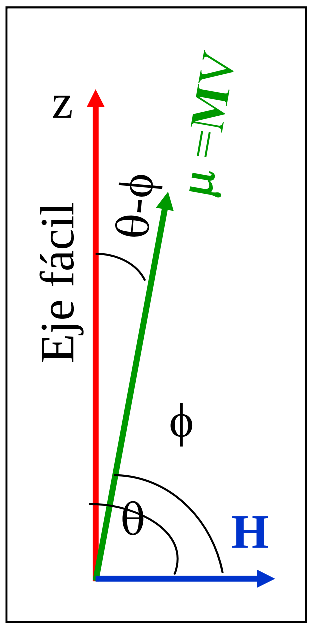
$$\theta = \pi / 2$$

$$E = E_K + E_H = KV\left[\sin^2(\phi - \pi/2) - 2h\cos\phi\right] = KV\left(\cos^2(\phi) - 2h\cos\phi\right)$$

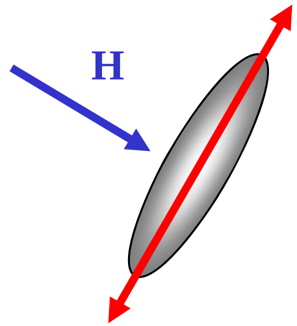
$$E = KV\cos\phi\left(\cos(\phi) - 2h\right)$$







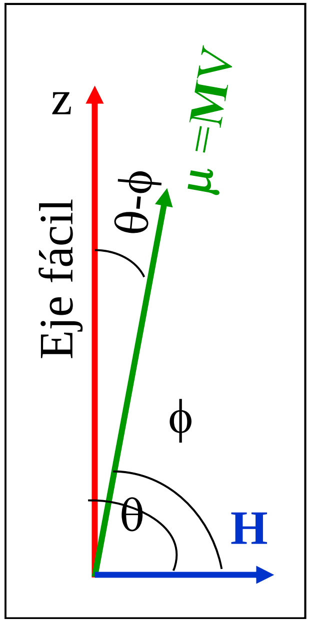
$$E = KV\cos\phi(\cos\phi - 2h)$$



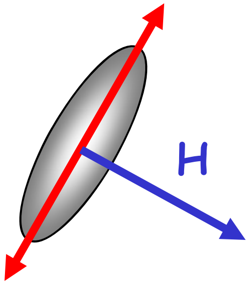
extremo $\Rightarrow \begin{cases} \sin\phi = 0 \Rightarrow \phi = \begin{cases} 0 \\ \pi \end{cases} \\ \cos\phi = h; |h| < 1 \end{cases}$

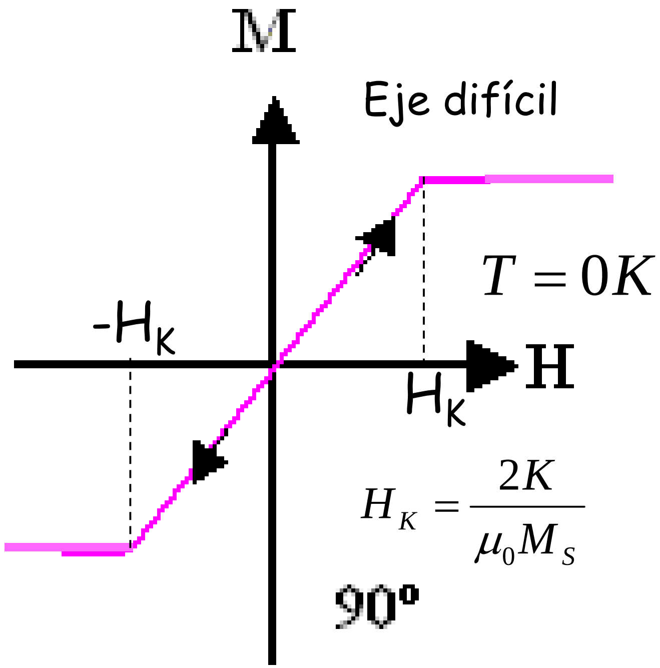
$$\frac{M_z}{M_S} = \cos\phi \qquad h = \frac{H}{H_K}$$

$$M_z = \frac{M_S}{H_K}H; \qquad |h| < 1$$



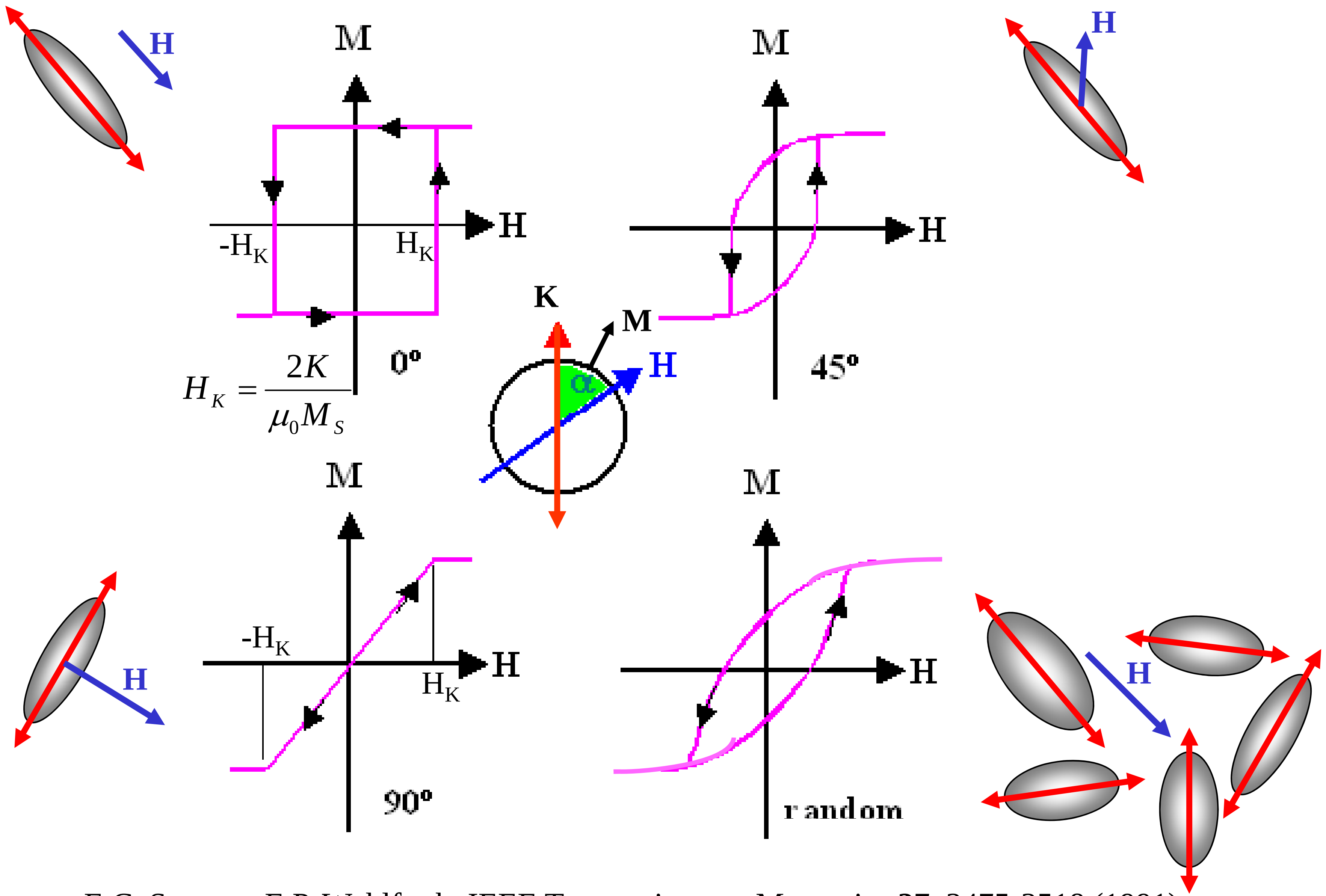
$$M_z = \frac{M_S}{H_K} H; \qquad |h| < 1$$





Partículas ferromagnéticas pequeñas – modelo de Stoner - Wohlfarth

régimen bloqueado → T = 0



E.C. Stoner y E.P. Wohlfarth, IEEE Transactions on Magnetics **27**, 3475-3518 (1991)

IBM Journal of Research and Development
Spintronics Volume 50, Number 1, 2006

Rapid-turnaround characterization methods for MRAM development

by D. W. Abraham,
P. L. Trouilloud,
and D. C. Worledge

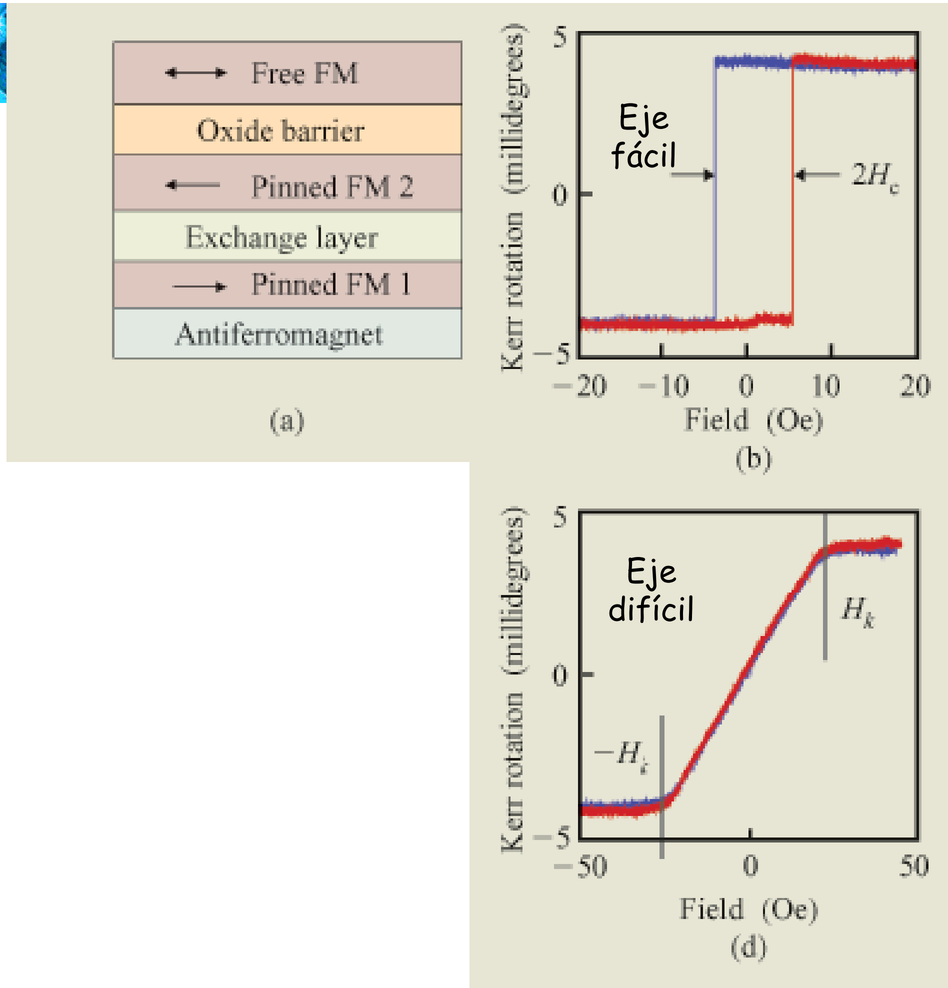


Figure 4

(a) Typical data for a Stoner–Wohlfarth stack. (a) Kerr easy-axis (EA) data taken at low field, showing the excellent low Néel offset and sharp hysteresis loop. (c) High-field EA Kerr magnetometry data showing the relative motion of the magnetization in the two ferromagnetic films, permitting direct measurement of pinning and interlayer coupling. (d) Hard-axis data revealing the film anisotropy.

A MECHANISM OF MAGNETIC HYSTERESIS IN HETEROGENEOUS ALLOYS

By E. C. STONER, F.R.S. and E. P. WOHLFARTH
Physics Department, University of Leeds

(*Received* 24 *July* 1947)

Vol. 240. A. 826 (Price 10*s*.) 74 [Published 4 May 1948

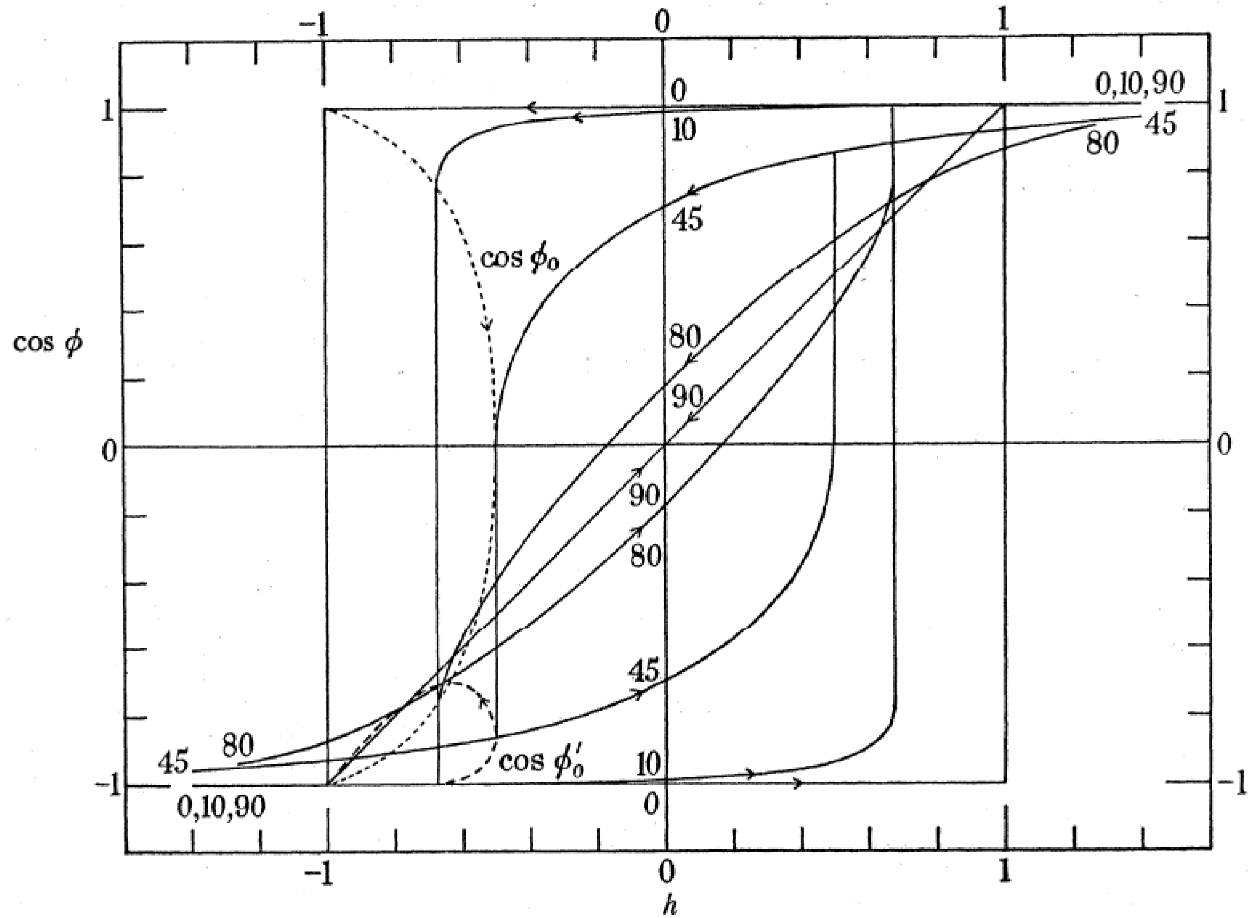


FIGURE 6. Magnetization curves for prolate spheroids. The resolved magnetization in the positive field direction is given by $I_0 \cos \phi$, where I_0 is the saturation magnetization. The field, H, is given by $H = (N_b - N_a) I_0 h$, where N_a and N_b are the demagnetization coefficients along the polar and equatorial axes. The angle, θ, between the polar axis and the direction of the field, is shown, in degrees, by the numbers on the curves. The dotted curves give $\cos \phi_0$ and $\cos \phi_0'$, where ϕ_0 and ϕ_0' are the angles made with the positive field direction by the magnetization vector at the beginning and end of the discontinuous change at the critical value, h_0, of the field.

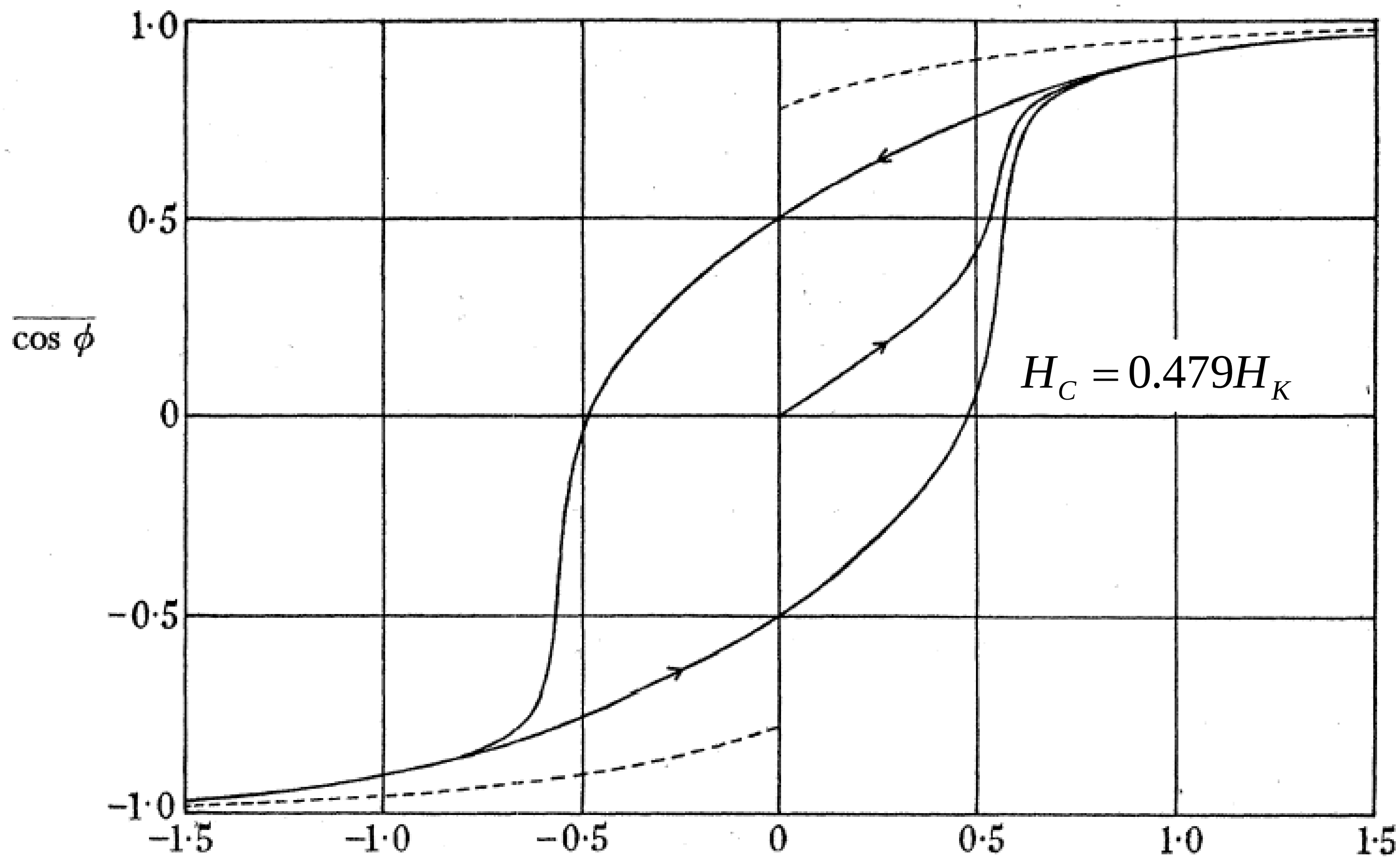
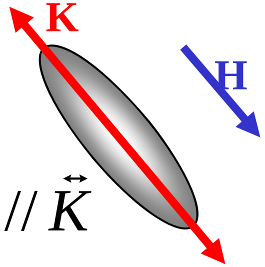


FIGURE 7. Magnetization curves for prolate (full curves) and oblate (broken curves) spheroids orientated at random. The curves refer to similar prolate (or oblate) spheroids orientated at random. $\overline{\cos\phi}$ is proportional to the mean resolved magnetization per spheroid in the positive field direction, or to the resultant magnetization in this direction of the assembly. $H = (|N_a - N_b|)\, I_0 h$.

Efectos Dinámicos ($T \neq 0$)

$$E = KV\left(\sin^2\phi - 2h\cos\phi\right)$$



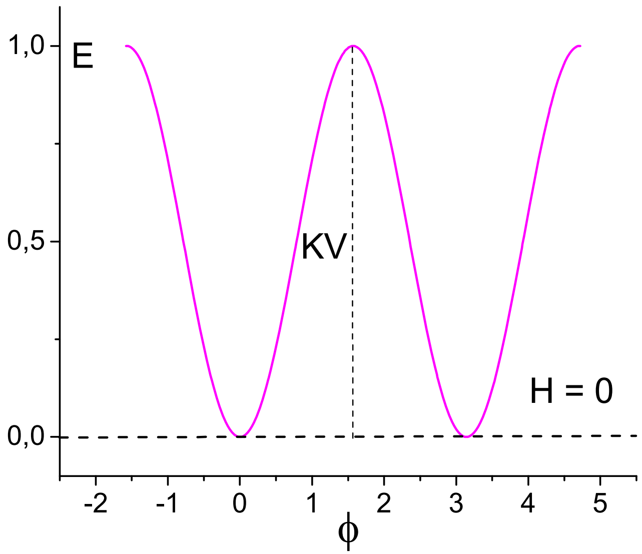
$$\vec{H} \,//\, \overleftrightarrow{K}$$

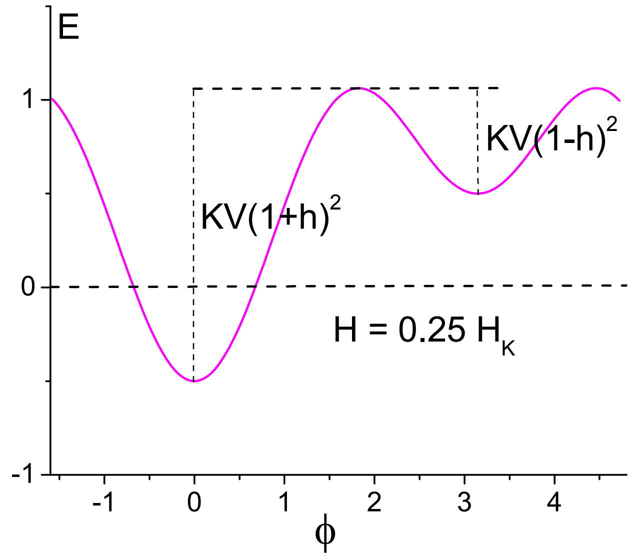
Valores de la función

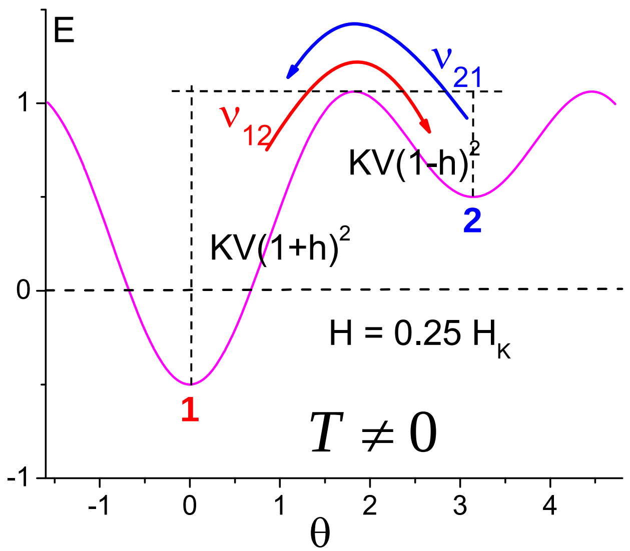
$E(\phi = 0) = -2hKV$	mín o máx
$E(\phi = \pi) = 2hKV$	mín o máx
$E(\cos\phi = -h) = KV(1 + h^2)$	máx o inflexión

$$h = \frac{H}{H_K}$$

$$H_K = \frac{2K}{\mu_0 M_S}$$







$$\nu_{ij} = c_0 e^{-\frac{\Delta E_{ij}}{kT}}$$

Frecuencia de saltos

$$\xrightarrow[T=0]{} \nu_{ij} = 0$$

$$\xrightarrow[T=\infty]{} \nu_{ij} = c_0$$

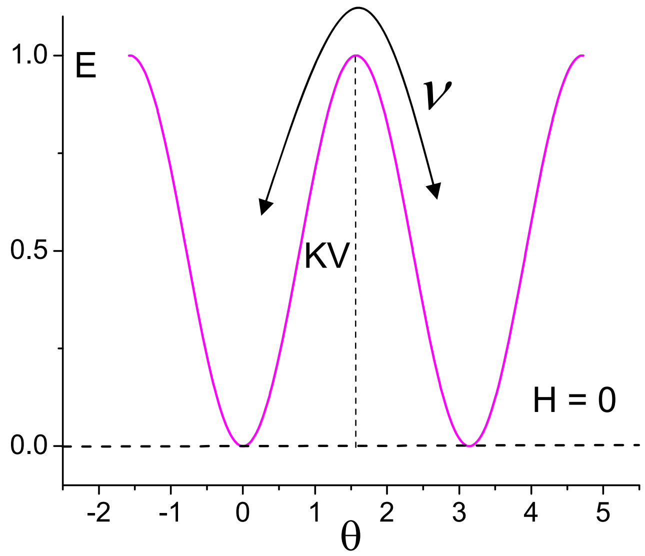
Frecuencia de intentos

$$\tau_{ij} = c_0^{-1} e^{\frac{\Delta E_{ij}}{kT}}$$

Tiempo de relajación

$$\tau_{ij} = c_0^{-1} e^{\frac{\Delta E_{ij}}{kT}}$$

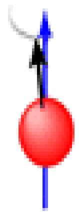
Tiempo de relajación



Para H = 0

$$\nu_{12} = \nu_{21} = \nu$$

$$\boxed{\nu = \nu_0 e^{-\frac{KV}{kT}}}$$



$$\boxed{\tau = \tau_0 e^{\frac{KV}{kT}}}$$

$$10^{-12}\, s \le \tau_0 \le 10^{-9}\, s$$

Estructura de τ_0

$$\boxed{\tau = \tau_0 e^{\frac{KV}{kT}}}$$

$$\tau_0 \approx cte$$

Modelo de Brown

$$\tau_0 \approx \frac{\mu}{2KV\gamma_0}\sqrt{\frac{\pi}{\alpha}} \qquad \alpha \approx \frac{KV}{kT}$$

$$M_S(T)$$

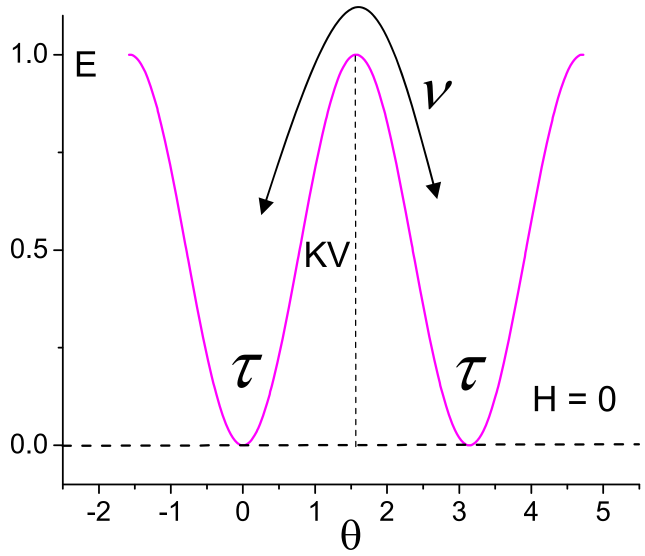
Ejemplo, usando $\tau_0 = 10^{-9}$ s

material	$K(J/m^3)$	R(nm)	τ(s)
Co	3.9×10^5	4.4	6×10^5
		3.6	0.1
Fe	4.7×10^4	14.0	1.5×10^5
		11.5	0.07

Comportamiento superparamagnético

$$\tau = \tau_0 e^{\frac{KV}{kT}}$$

Tiempo Experimental vs Tiempo de Relajación

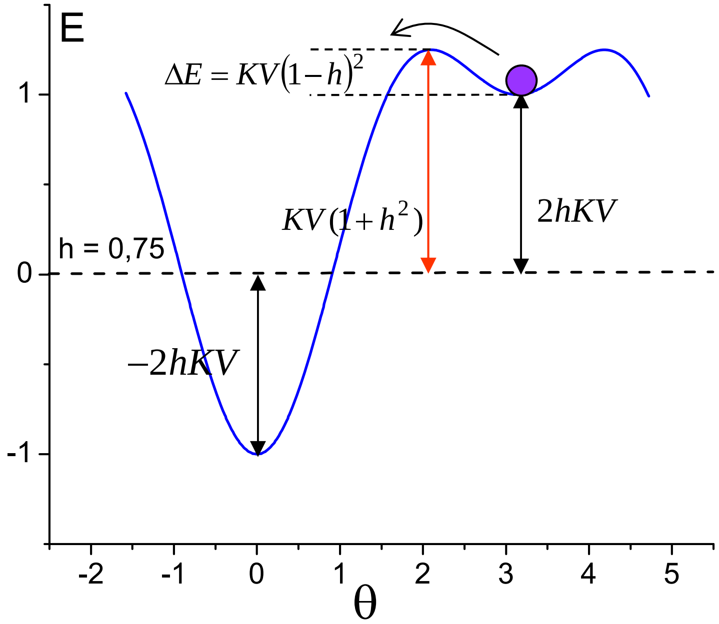


Técnica	τ_{exp}
Mössbauer ^{57}Fe, ^{119m}Sn	$\approx 10^{-8}$ s
Susceptibilidad *ac*	10^{-4} – 1 s
Susceptibilidad *ac hf*	desde 10^{-6} s
Magnetización *dc*	0.1 – 100 s

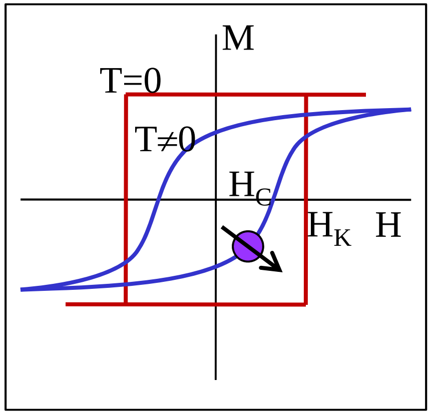
$\tau_{exp} < \tau \Leftrightarrow T < T_B$	Sistema bloqueado	Patrón estático	Histéresis, desdoblamiento Zeeman (EM)
$\tau_{exp} > \tau \Leftrightarrow T > T_B$	Sistema desbloqueado	Patrón dinámico	Equilibrio, patrón super-paramagnético (EM)

Dependencia del campo coercitivo con la temperatura

$$h = H / H_K = \frac{\mu_0 M_S H}{2K}$$



$$H_C = H_K = \frac{2K}{\mu_0 M_S}$$



$$\tau_{21} = \tau_0 e^{\frac{KV(1-h)^2}{kT}}$$

a T ≠ 0 K la inversión de M se producirá cuando $\tau_{21} \approx \tau_{exp}$

Dependencia del campo coercitivo con la temperatura

$$KV(1-h)^2 = kT\ln(\tau_{\exp}/\tau_0)$$

$\tau_{\exp} = 10^2 s$ $SQUID$

$\tau_{\exp} = 10^{-8} s$ $Möss$

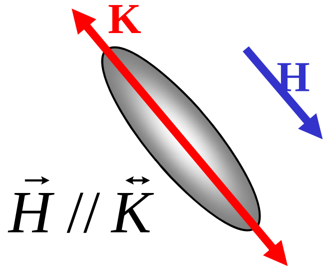
$$\ln(\tau_{\exp}/\tau_0) \approx \begin{cases} 27.6\,kT & SQUID \\ 4.6\,kT & Möss \end{cases}$$

$$\left.\begin{array}{l} KV(1-h)^2 \approx 27.6\,kT \\ KV(1-h)^2 \approx 4.6\,kT \end{array}\right. \xrightarrow[h<<1]{} \begin{array}{ll} KV \approx 27.6\,kT & SQUID \\ KV \approx 4.6\,kT & Möss \end{array}$$

Dependencia del campo coercitivo con la temperatura

$$KV(1-h)^2 \approx kT\ln\left(\tau_{\exp}/\tau_0\right)$$

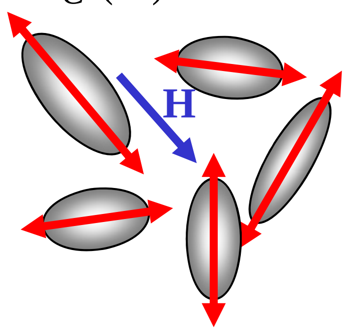
$$h \approx 1-\sqrt{\frac{kT}{KV}\ln\left(\tau_{\exp}/\tau_0\right)}$$

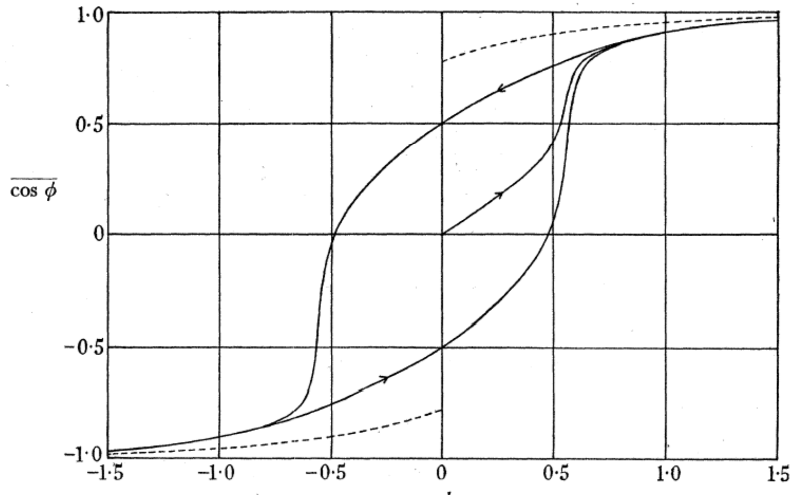


$$H_C(T) \approx H_K\left(1-\sqrt{\frac{kT}{KV}\ln\left(\tau_{\exp}/\tau_0\right)}\right)$$

$$H_C(0) = 0.479 H_K$$

$$H_C(T) \approx 0.48 H_K\left(1-\sqrt{\frac{kT}{KV}\ln\left(\tau_{\exp}/\tau_0\right)}\right)$$

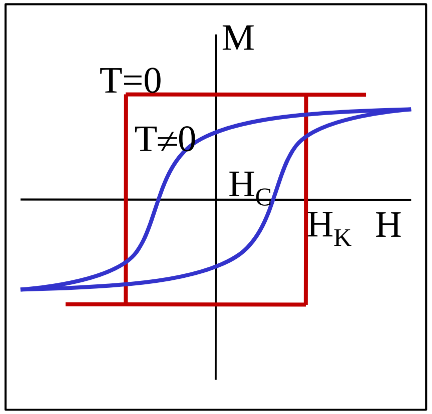




Temperature Dependent Magnetic Properties of Barium-Ferrite Thin-Film Recording Media

Yingjian Chen, *Member, IEEE,* and Mark H. Kryder, *Fellow, IEEE*

IEEE TRANSACTIONS ON MAGNETICS, VOL. 34, NO. 3, MAY 1998



$$H_c(t') = H_k\left\{1 - \left[\frac{k_B T}{K_u V_{sw}} \ln\left(\frac{At'}{0.693}\right)\right]^n\right\}$$

the easy axis orientation. In a system with uniaxially aligned easy axes, n is 1/2 [29], and in a system with random easy axis orientations n is 2/3 [30]. The fitting parameters V_{sw}

[29] M. P. Sharrock and J. T. McKinney, *IEEE Trans. Magn.*, vol. MAG-17, p. 3020, 1981.
[30] R. H. Victora, "Predicted time dependence of the switching field for magnetic materials," *Phys. Rev. Lett.*, vol. 63, pp. 457–460, 1989.

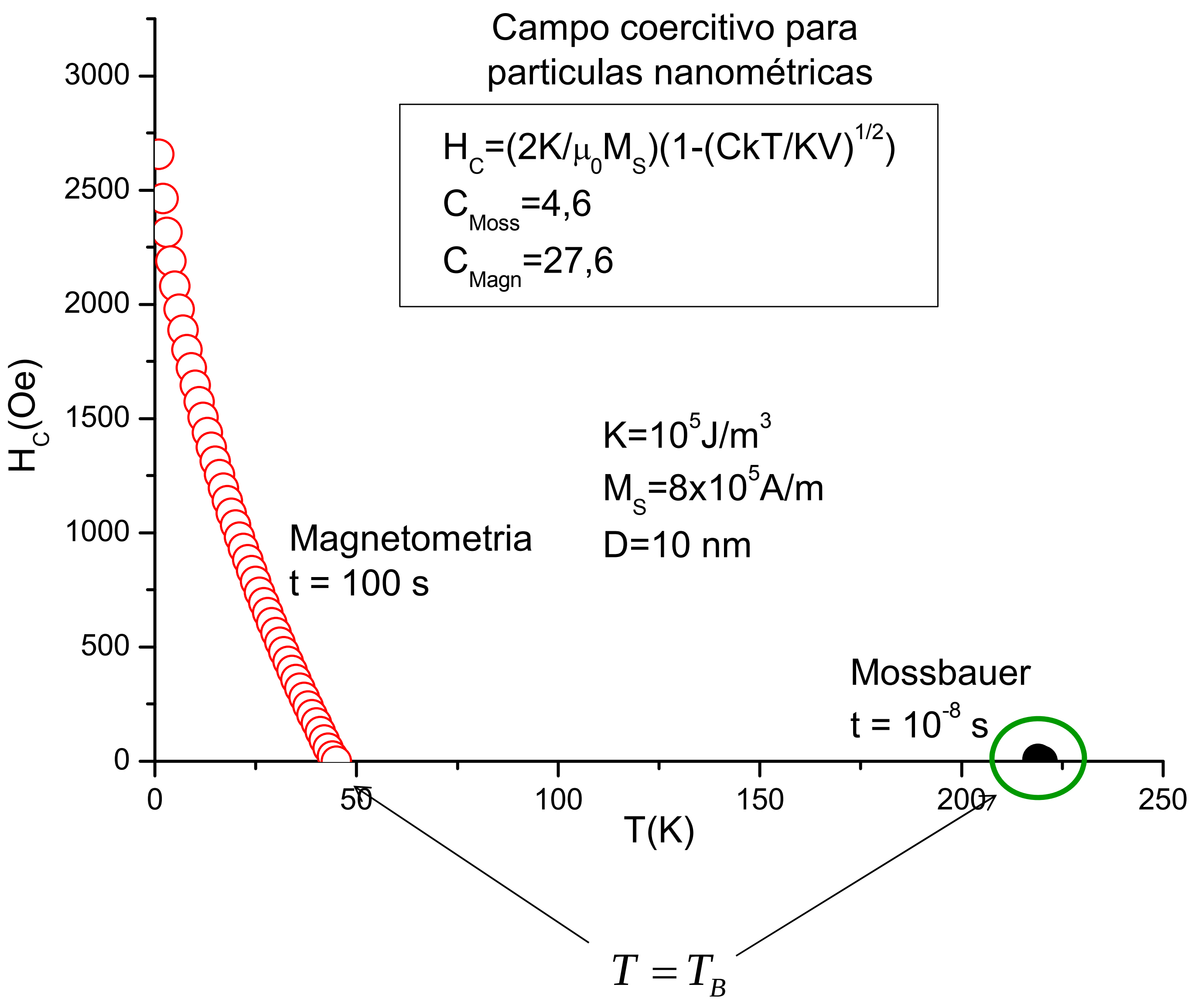
Uso extendido de la expresión

$$H_C = \alpha \frac{2K}{M_S}\left[1 - \left(\frac{T}{T_B}\right)^{1/2}\right]$$

Interacciones magnéticas en nanotubos ferromagnéticos de LaCaMnO y LaSrMnO,

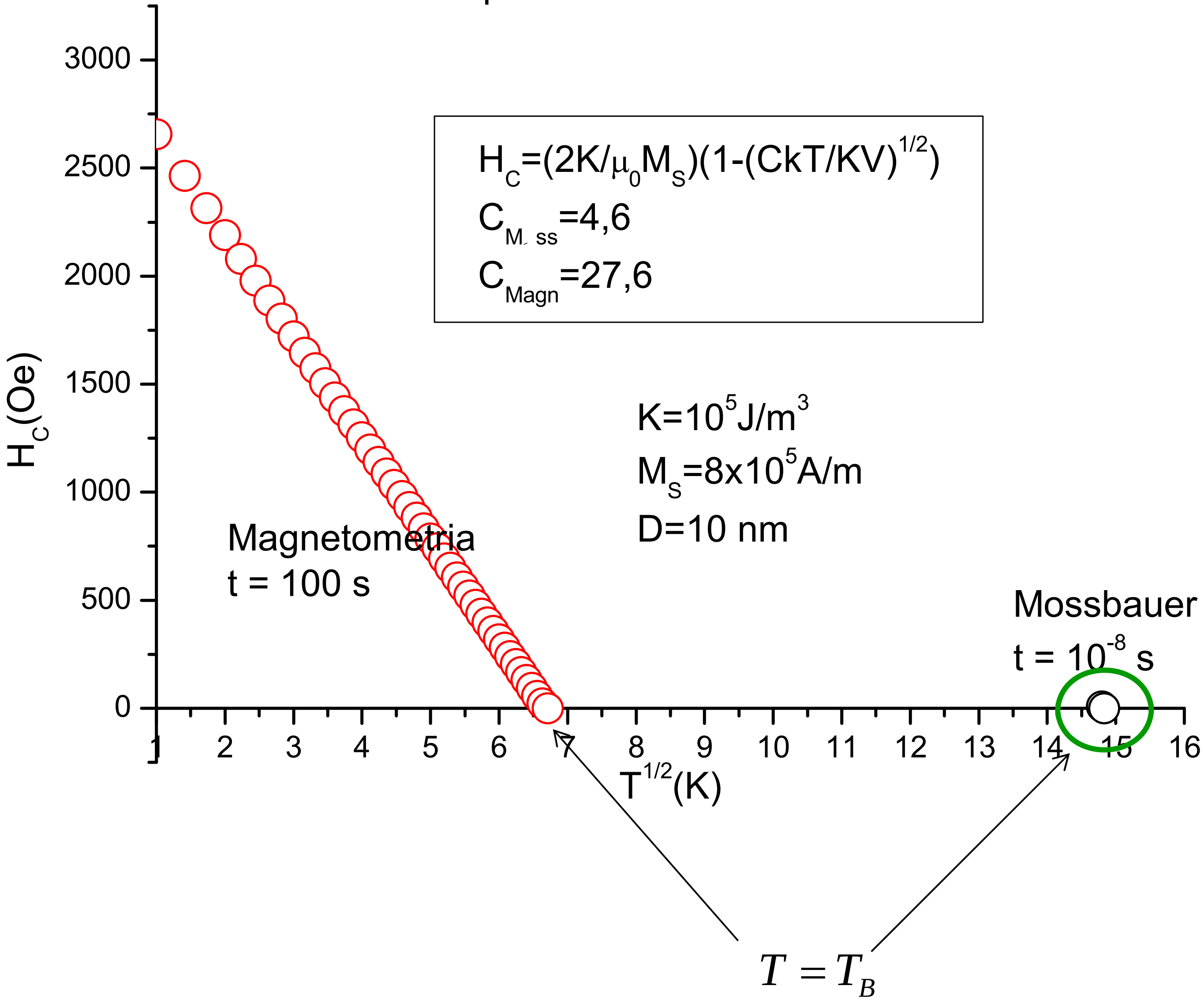
J.Curiale et al., AFA 2006

Marina Tortarola, Tesis, IB, 2008



Campo coercitivo para
particulas nanométricas
$H_C=(2K/\mu_0 M_S)(1-(CkT/KV)^{1/2})$
$C_{Moss}=4,6$
$C_{Magn}=27,6$
$K=10^5 J/m^3$
$M_S=8x10^5 A/m$
D=10 nm
Magnetometria
t = 100 s
Mossbauer
$t = 10^{-8}$ s
H_C(Oe)
T(K)
0
500
1000
1500
2000
2500
3000
50
100
150
200
250
$T = T_B$

Campo coercitivo para
particulas nanométricas



Fin módulo